

The Logarithmic Schrödinger-Langevin Equation and Quantum Fluids

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Logarithmic Schrödinger-Langevin Equation

$$\boxed{i\partial_t\psi + \frac{1}{2}\Delta\psi = \lambda\psi \log(|\psi|^2) + \frac{1}{2i}\mu\psi \log\left(\frac{\psi}{\psi^*}\right)}, \quad (1)$$

for $t \in \mathbb{R}_+$ and $x \in \mathbb{R}^d$.

\Rightarrow physics : $\mu > 0$. Influence of $\lambda > 0$ and $\lambda < 0$?

- well-posedness?
- local/global existence? long time behavior?
- stationnary solutions? stability?
- numerics?

- 1 Schrödinger-Langevin Equation
 - Gaussian solutions and related ODE
 - $\lambda < 0$ case
 - $\lambda > 0$ case

- 2 Quantum Hydrodynamics Systems
 - Isothermal Euler-Langevin-Korteweg Equation
 - Long-time Behavior
 - Splitting Method

Logarithmic Schrödinger Equations and Quantum Fluids

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Gaussian solutions and related ODE

Plug

$$\psi(t, x) = b(t)e^{-\frac{1}{2}a(t)x^2}$$

into (1), $a, b \in \mathbb{C}$. After calculations, you get

$$\ddot{r} = \frac{\alpha_0^2}{r^3} + \frac{2\lambda\alpha_0}{r} - \mu\dot{r}, \quad (2)$$

where

$$a(t) = \frac{\alpha_0}{r(t)^2} - i\frac{\dot{r}(t)}{r(t)},$$

$$|b(t)| = |b_0| \exp\left(\frac{1}{2} \operatorname{Im} A(t)\right),$$

with $A(t) = \int_0^t a(s)ds$, $\alpha_0 = \operatorname{Re} a_0 \geq 0$ and $r \in \mathbb{R}$.

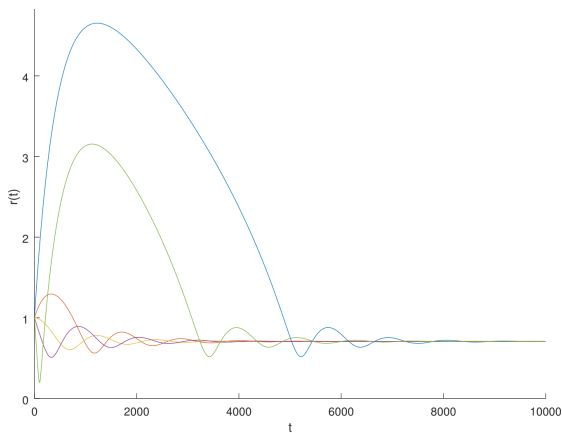
$\lambda < 0$ case

Figure – Dynamic of equation (2)

$\lambda < 0$ case

Proposition

$$r(t) \xrightarrow{t \rightarrow +\infty} \sqrt{\frac{\alpha_0}{-2\lambda}}$$
$$\dot{r}(t) \xrightarrow{t \rightarrow +\infty} 0$$

Corollary

As $a(t) \rightarrow -2\lambda$, every Gaussian solution ψ tends to a Gaussian

$$\phi(x) = C_{\lambda, \alpha_0, \beta_0} e^{\lambda x^2},$$

uniquely determined by the IC α_0 and β_0 .

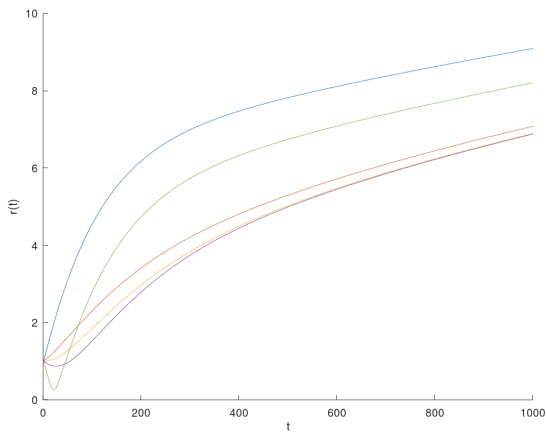
$\lambda > 0$ case

Figure – Dynamic of equation (2)

$\lambda > 0$ case

Proposition

$$r(t) \sim 2\sqrt{\frac{\lambda\alpha_0}{\mu}}t$$

$$\dot{r}(t) \sim \sqrt{\frac{\lambda\alpha_0}{\mu t}}$$

Corollary

$$\|\psi(t)\|_{L^\infty} \sim C_{\mu,\lambda,\alpha_0} t^{-\frac{1}{4}}$$

$$\|\nabla\psi(t)\|_{L^2}^2 \sim C_{\mu,\lambda,\alpha_0,|b_0|} \frac{1}{\sqrt{t}}$$

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Isothermal Euler-Langevin-Korteweg Equation

We recall :

$$i\hbar\partial_t\psi + \frac{\hbar^2}{2}\Delta\psi = \lambda\psi\log(|\psi|^2) + \frac{\hbar}{2i}\mu\psi\log\left(\frac{\psi}{\psi^*}\right)$$

Madelung transform $\psi = \sqrt{\rho}e^{iS/\hbar}$, and define $J = \rho\nabla S/\hbar$, then :

$$\partial_t\rho + \operatorname{div}J = 0 \quad (3)$$

$$\partial_t J + \operatorname{div}\left(\frac{J \otimes J}{\rho}\right) + \lambda\nabla\rho + \mu J = \frac{\hbar^2}{2}\rho\nabla\left(\frac{\Delta\sqrt{\rho}}{\sqrt{\rho}}\right) \quad (4)$$

Remark

Pathological 3rd order quantum potential $\rho\nabla\left(\frac{\Delta\sqrt{\rho}}{\sqrt{\rho}}\right)$.

Long-time Behavior

Proposition

$\lambda < 0$ case : the density ρ of every solution of (3)-(4) converges to a stationary Gaussian $\phi(x) = Ce^{\lambda|x-x_\infty|^2}$ weakly in $L^1(\mathbb{R}^d)$.

Proposition

$\lambda > 0$ case : define

$$\rho(t, x) = \frac{1}{\tau(t)^d} R\left(t, \frac{x}{\tau(t)}\right)$$

where τ is the unique solution of the nonlinear ODE : $\ddot{\tau} = \frac{2\lambda}{\tau} - \mu\dot{\tau}$.
Then

$$R(t, x) \rightharpoonup Ce^{-\frac{|x|^2}{2}} \quad \text{weakly in } L^1(\mathbb{R}^d).$$

Splitting Method

$$i\partial_t\psi + \frac{1}{2}\Delta\psi = \lambda\psi \log(|\psi|^2) + \frac{1}{2i}\mu\psi \log\left(\frac{\psi}{\psi^*}\right)$$

We solve :

- $\partial_t\psi = -\frac{1}{2}i\Delta\psi$ by FFT,
- $\partial_t\psi = -i\lambda\psi \log(|\psi|^2 + \varepsilon)$ by the explicit solution

$$\psi(t + \Delta t, \cdot) = \psi(t, \cdot)e^{-i\lambda\Delta t \log(|\psi(t, \cdot)|^2 + \varepsilon)},$$

- and $\partial_t\psi = -\frac{1}{2}\mu\psi \log\left(\frac{\psi}{\psi^*}\right)$ by an explicit solution

$$\psi(t + \Delta t, \cdot) = a(t, \cdot)e^{i\theta(t, \cdot)}e^{-\mu\Delta t},$$

where we decompose $\psi(t, \cdot) = a(t, \cdot)e^{i\theta(t, \cdot)}$.

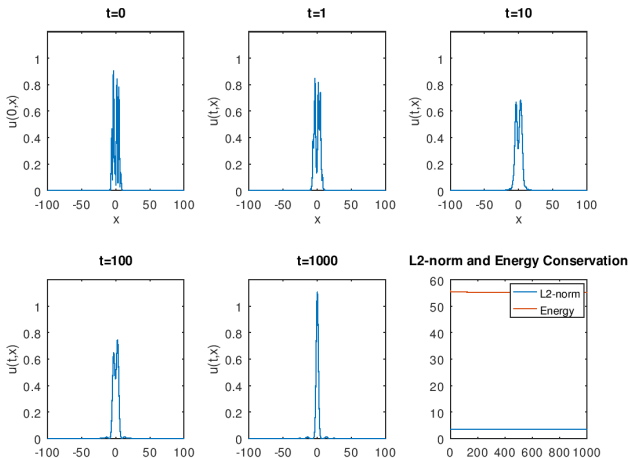
$\lambda < 0$ case

Figure – Solution of equation (1) with initial datum ψ_0 in the focusing case ($\lambda = -0.1$, $\mu = 1$).

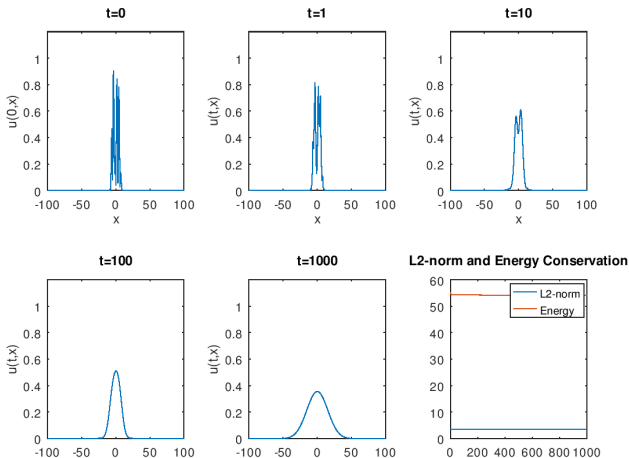
$\lambda > 0$ case

Figure – Solution of equation (1) with initial datum ψ_0 in the defocusing case ($\lambda = 0.1$, $\mu = 1$).

References



Rémi Carles and Isabelle Gallagher.

Universal dynamics for the defocusing logarithmic Schrödinger equation.

Duke Math. J., 167(9) :1761–1801, 2018.



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Dynamics of the Schrödinger-Langevin equation.

Nonlinearity, 34(4) :1943–1974, 2021.

Thank you for your attention.