

# Hyperbolic Quadrature Method of Moments for the one-dimensional kinetic equation

**Frédérique LAURENT-NÈGRE**

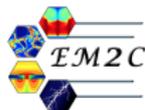
Laboratoire EM2C - CNRS - CentraleSupélec - Université Paris-Saclay  
Fédération de Mathématiques de CentraleSupélec - CNRS

**Rodney O. Fox**

Department of Chemical and Biological Engineering, Iowa State University, USA

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# Considered kinetic model

Kinetic model on  $f(t, \mathbf{x}, \mathbf{v})$

$$\underbrace{\partial_t f + \partial_{\mathbf{x}} \cdot (\mathbf{v} f)}_{\text{physical transport}} = \underbrace{S(f)}_{\text{Source terms}}$$

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Gas dynamics

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Transition regime :  $0.01 < Kn < 10$

Far from the Maxwellian equilibrium

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- The kinetic model is too costly to solve with direct methods of Monte-Carlo type
- Moments  $\int_{\mathbb{R}} \mathbf{v}^k f(t, \mathbf{x}, \mathbf{v}) d\mathbf{v}$  of order  $k$  smaller than 1 or 2 are not enough to represent the distribution.

# Moment method

## Principle of the method

Write equations on a finite set of moments  $\mathbf{m}_N = (m_0, m_1, \dots, m_N)^t$ :

$$\partial_t m_k + \partial_x m_{k+1} = S_k, \quad k = 0, 1, \dots, N \quad (1)$$

**Closure:** express  $m_{N+1}$  (and eventually the source terms  $S_k$ ) as a function of  $\mathbf{m}_N$ .

### Issues:

- $(m_0, m_1, \dots, m_N, m_{N+1})^t$  is realizable
- The system (1) is globally hyperbolic
- Capture equilibrium state

### Strategy

- Solve the Hamburger truncated moment problem:

$$\text{find a positive measure } \mu \text{ such that } \mathbf{m}_N = \int_{\mathbb{R}} (1, v, \dots, v^N)^t d\mu(v).$$

and set  $m_{N+1} = \int_{\mathbb{R}} v^{N+1} d\mu(v)$

- Give directly  $m_{N+1}$

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## Examples of closure in the literature

- Grad closure [Grad, 1949]
  - hyperbolic only around the moments of the maxwellian distribution
- Entropy maximization [Levermore, 1996, Müller and Ruggeri, 1998]
  - high computational cost - not valid on the entire realizability domain
- Quadrature method of moment [McGraw, 1997, Fox, 2008]
  - weakly hyperbolic

# Moment space

## Definition

The  $n^{\text{th}}$ -moment space  $\mathcal{M}_n$  is defined by

$$\mathcal{M}_n = \left\{ \mathbf{m} \in \mathbb{R}^{n+1} \mid \exists \mu \in \mathcal{M}_+(\mathbb{R}), \quad \mathbf{m} = \int_{\mathbb{R}} (1, v, \dots, v^n)^t d\mu(v) \right\}$$

If  $\mathbf{m}$  belongs to  $\mathcal{M}_n$ , then it is said to be realizable.

If  $\mathbf{m}$  belongs to the interior  $\text{Int } \mathcal{M}_n$  of  $\mathcal{M}_n$ , it is said to be strictly realizable.

Characterized by the non-negativity of the Hankel determinants:  $n \geq 0$

$$\underline{H}_{2n} = \begin{vmatrix} m_0 & \dots & m_n \\ \vdots & & \vdots \\ m_n & \dots & m_{2n} \end{vmatrix},$$

## Theorem

$\mathbf{m}_N = (m_0, m_1, \dots, m_N)^t$  strictly realizable  $\Leftrightarrow \underline{H}_{2k} > 0, \quad k \in \{0, 1, \dots, \lfloor \frac{N}{2} \rfloor\}$

$\mathbf{m}_N \in \partial \mathcal{M}_N \cap \mathcal{M}_n \Rightarrow \underline{H}_0 > 0, \dots, \underline{H}_{2k-2} > 0, \underline{H}_{2k} = 0, \dots, \underline{H}_{2\lfloor \frac{N}{2} \rfloor} = 0, k \leq \lfloor \frac{N}{2} \rfloor$ .

In the latter case, the only corresponding measure is a sum of  $k$  weighted Dirac delta functions.

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Characterized by the non-negativity of the **Hankel determinants**:  $n \geq 0$

$$\underline{H}_{2n} = \begin{vmatrix} m_0 & \dots & m_n \\ \vdots & & \vdots \\ m_n & \dots & m_{2n} \end{vmatrix},$$

First constraints for the strict realizability:

$$\begin{aligned} m_0 &> 0 & m_2 &> \frac{m_1^2}{m_0} \\ m_4 &> \frac{m_0 m_3^2 - 2m_1 m_2 m_3 + m_2^3}{m_2 m_0 - m_1^2} \\ && \dots & \end{aligned}$$

# Characteristic polynomial

## System on moments

Equations on  $\mathbf{m}_N = (m_0, m_1, \dots, m_N)^t$ :

$$\partial_t \mathbf{m}_N + \partial_x F(\mathbf{m}_N) = \bar{S}$$

## Characteristic Polynomial

Jacobian matrix

$$\frac{DF(\mathbf{m}_N)}{D\mathbf{m}_N} = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & & 0 \\ \vdots & & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & 1 \\ \frac{\partial m_{N+1}}{\partial m_0} & \frac{\partial m_{N+1}}{\partial m_1} & \frac{\partial m_{N+1}}{\partial m_2} & \dots & \frac{\partial m_{N+1}}{\partial m_N} \end{pmatrix}.$$

Characteristic polynomial

$$\bar{P}_{N+1}(X) = X^{N+1} - \sum_{i=0}^N \frac{\partial m_{N+1}}{\partial m_i} X^i$$

# Moments - Central moments - Standardized moments

**moments:**

$$m_k = \int_{\mathbb{R}} v^k f(v) dv$$

**central moments:** with  $\rho = m_0$ ,  $u = \frac{m_1}{m_0}$  and  $f^c(c) = \frac{1}{\rho} f(c + u)$

$$C_k = \frac{1}{\rho} \int_{\mathbb{R}} (v - u)^k f(v) dv = \int_{\mathbb{R}} c^k f^c(c) dc$$

so that  $C_0 = 1$  and  $C_1 = 0$ .

**standardized moments:** with  $\sigma = \sqrt{C_2}$ ,  $f^s(s) = \frac{\sigma}{\rho} f(u + \sigma s)$

$$S_k = \frac{1}{m_0} \int_{\mathbb{R}} \left( \frac{v - u}{\sqrt{C_2}} \right)^k f(v) dv = \int_{\mathbb{R}} s^k f^s(s) ds$$

so that  $S_0 = 1$ ,  $S_1 = 0$  and  $S_2 = 1$ .

**link:**

$$C_k = \sum_{i=0}^k \binom{k}{i} \left( -\frac{m_1}{m_0} \right)^{k-i} m_i, \quad m_k = \rho \left( \sum_{i=2}^k \binom{k}{i} u^{k-i} C_i + u^k \right), \quad S_k = \frac{C_k}{(C_2)^{k/2}}.$$

# Property of the characteristic polynomial

$\mathbf{m}_N = (m_0, m_1, \dots, m_N)^t$  be a realizable moment vector such that  $m_0 > 0$  and  $C_2 > 0$ .

linear functional  $\langle \cdot \rangle_{\mathbf{m}_N}$  on the space  $\mathbb{R}[X]_N$

$$\langle X^k \rangle_{\mathbf{m}_N} = m_k, \quad \text{for } k \in \{0, 1, \dots, N\}.$$

linear functional  $\langle \cdot \rangle_{\mathbf{S}_N}$  associated with the standardized moments  $\mathbf{S}_N = (S_0, \dots, S_N)^t$ :

$$\langle X^k \rangle := \langle X^k \rangle_{\mathbf{S}_N} = S_k, \quad \text{for } k \in \{0, 1, \dots, N\}.$$

## Property of the scaled characteristic polynomial

Let us assume that the function  $S_{N+1}$  does not depend on  $(m_0, u, C_2)$ , i.e.,  $S_{N+1}(S_3, \dots, S_N)$ . Then, the following polynomial

$$P_{N+1}(X) := \bar{P}_{N+1} \left( u + C_2^{1/2} X \right) C_2^{-(N+1)/2}$$

only depends on  $(S_3, \dots, S_N)$ , and

$$\langle P_{N+1} \rangle = 0, \quad \langle P'_{N+1} \rangle = 0, \quad \langle X P'_{N+1} \rangle = 0.$$

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# QMOM: Principles of the method

From a strictly realizable moment vector  $\mathbf{m}_{2n-1}$

## Reconstruction

reconstruct the discret measure  $\mu = \sum_{i=1}^n w_i \delta_{u_i}$

in such a way that

$$\sum_{i=1}^n w_i u_i^k = m_k \quad k = 0, 1, \dots, 2n-1$$

## Closure

$$m_{2n} = \int_{\mathbb{R}} v^{2n} d\mu = \sum_{i=1}^n w_i u_i^{2n}$$

It is the minimal value for this moment

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From the standardized moments  $\mathbf{S}_{2n-1}$ ,  
with  $\rho = m_0$ ,  $u = m_1/m_0$ ,  $\sigma = \sqrt{C_2}$

## Reconstruction

reconstruct the discret measure

$\mu = \sum_{i=1}^n \rho \omega_i \delta_{u+\sigma c_i}$  in such a way that

$$\sum_{i=1}^n \omega_i c_i^k = S_k \quad k = 0, 1, \dots, 2n-1$$

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## Remarks

- The reconstruction  $\mu$  is the only one possible for the moment vector  $\mathbf{m}_{2n}$ .
- $\mathbf{m}_{2n}$  is at the boundary of the moment space:  $H_{2n} = 0$

# QMOM: Computation of the weights and abscissas

$\mathbf{m}_{2n-1}$ : strictly realizable moment vector

## Orthogonal polynomials

Family  $(Q_k)_{k=0,\dots,n}$  of monic orthogonal polynomials for the scalar product  $(p, q) \mapsto \langle pq \rangle$  of  $\mathbb{R}_n[X]$ .

$$Q_{k+1}(X) = (X - a_k)Q_k(X) - b_k Q_{k-1}(X)$$

with  $Q_{-1} = 0$  and  $Q_0 = 1$ .

The recurrence coefficients  $a_k$  and  $b_k$  can be found from the standardized moments using the Chebyshev algorithm [Chebyshev, 1859, Wheeler, 1974, Gautschi, 2004]

$$a_k = \frac{\langle XQ_k^2 \rangle}{\langle Q_k^2 \rangle}, \quad b_k = \frac{\langle Q_k^2 \rangle}{\langle Q_{k-1}^2 \rangle} = \frac{H_{2k}H_{2k-4}}{H_{2k-2}^2}.$$

example

$$\begin{aligned} a_0 &= 0, & a_1 &= S_3, & a_2 &= \frac{S_5 - S_3(2 + S_3^2 + 2H_4)}{H_4} \\ b_0 &= 1, & b_1 &= 1, & b_2 &= H_4, & b_3 &= H_6/H_4^2 \end{aligned}$$

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The abscissas  $c_i$  are the zeros of  $Q_n$  and also the eigenvalues of the Jacobi matrix

$$\mathbf{J}_n = \begin{pmatrix} a_0 & \sqrt{b_1} & & & \\ \sqrt{b_1} & a_1 & & & \\ & & \sqrt{b_2} & & \\ & & \ddots & \ddots & \\ & & & \sqrt{b_{n-2}} & a_{n-2} & \sqrt{b_{n-1}} \\ & & & & \sqrt{b_{n-1}} & a_{n-1} \end{pmatrix}$$

# Hyperbolicity of the QMOM method

## Theorem

*The QMOM closure  $b_n = 0$  induces the following characteristic polynomial  $P_{2n} = Q_n^2$  and the system is only weakly hyperbolic.*

proof [Chalons et al., 2012, Huang et al., 2020]

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# First version of the HyQMOM closure [Fox et al., 2018]

Extension of QMOM, adding one moment and one abscissa for the reconstruction [Fox et al., 2018]

## three-node HyQMOM

reconstruction with an additional fixed abscissa  $\mu = w_0 \delta_u + \sum_{i=1}^2 w_i \delta_{u_i}$  in such a way that

$$w_0 u^k + \sum_{i=1}^n w_i u_i^k = m_k \quad k = 0, 1, \dots, 4$$

## Closure

in term of the standardized moments:  $S_5 = S_3(2S_4 - S_3^2)$

## Theorem (Hyperbolicity)

*Assuming that the vector  $\mathbf{m}_4$  is strictly realizable, then system with the three-node HyQMOM closure is hyperbolic.*

## Problem

- The generalization to a larger number of moment is not easy
- The eigenvalues of the problem do not tend to the ones of QMOM when  $\underline{H}_4 \rightarrow 0$

# New HyQMOM closure [Fox and Laurent, 2021]

## Idea:

- Instead of looking at a reconstruction or at a closure on  $\mathcal{S}_{2n+1}$ , one looks at  $a_n$ .
- Have a reduced characteristic polynomial on the form

$$P_{2n+1} = Q_n [(X - \alpha_n)Q_n - \beta_n Q_{n-1}]$$

such that  $\beta_n$  tends to zero when  $H_{2n} \rightarrow 0$ .

## Theorem

For all  $n = 1, 2, \dots$ ; let the monic polynomial  $P_{2n+1}$  be given by

$$P_{2n+1} = Q_n [(X - \alpha_n)Q_n - \beta_n Q_{n-1}] \quad \alpha_n, \beta_n \in \mathbb{R}$$

Then, the following statements are equivalent:

- $\langle P_{2n+1} \rangle = 0$ ,  $\langle P'_{2n+1} \rangle = 0$  and  $\langle XP'_{2n+1} \rangle = 0$ .
- $\alpha_n = a_n = \frac{1}{n} \sum_{k=0}^{n-1} a_k$  and  $\beta_n = \frac{2n+1}{n} b_n$ .

# New HyQMOM closure [Fox and Laurent, 2021]

## Theorem

For all  $n = 1, 2, \dots, 9$ ; the scaled characteristic polynomial can be written as

$$P_{2n+1} = Q_n [(X - \alpha_n)Q_n - \beta_n Q_{n-1}]$$

if and only if the closure on  $S_{2n+1}$ , defined through the coefficient  $a_n$ , and the coefficients  $\alpha_n$  and  $\beta_n$  are related to the recurrence coefficients  $a_k$  and  $b_k$  by

$$a_n = \alpha_n = \frac{1}{n} \sum_{k=0}^{n-1} a_k, \quad \beta_n = \frac{2n+1}{n} b_n.$$

Proof using formal computation with matlab symbolic:

from the  $a_k$  and  $b_k$ ,  $k = 0, \dots, n-1$  (with  $a_0 = 0$ ,  $a_1 = 1$ ,  $b_0 = 1$ )

- ① set the closure  $a_n = \frac{1}{n} \sum_{k=0}^{n-1} a_k$
- ② compute the Standardized moments  $S_{2n+1}$  with the reverse Chebyshev algorithm
- ③ compute the coefficients  $c_k$  of  $P_{2n+1}$
- ④ compute the polynomials  $Q_k$ ,  $k = 0, 1, \dots, n$
- ⑤ compute  $P_{2n+1} - Q_n \left[ (X - a_n)Q_n - \frac{2n+1}{n} b_n Q_{n-1} \right]$

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## Examples

- $n = 1$ :  $S_3 = 0$  (as for the Maxwellian reconstruction)
- $n = 2$ :  $S_5 = \frac{1}{2} S_3 (5S_4 - 3S_3^2 - 1)$  (different from the previous version:  
 $S_5 = S_3 (2S_4 - S_3^2)$ )



# Hyperbolicity - Eigenvalues

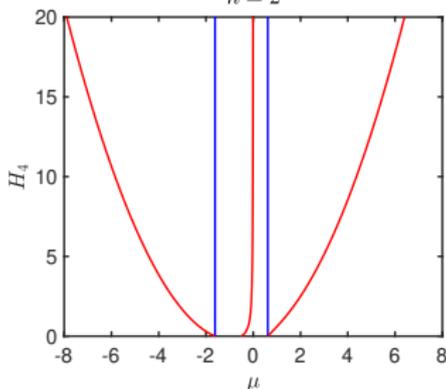
## Theorem

When  $\beta_n > 0$ , the  $n + 1$  roots of  $R_{n+1} = (X - \alpha_n)Q_n - \beta_n Q_{n-1}$  are real-valued and bound and separate the  $n$  roots of  $Q_n$ .

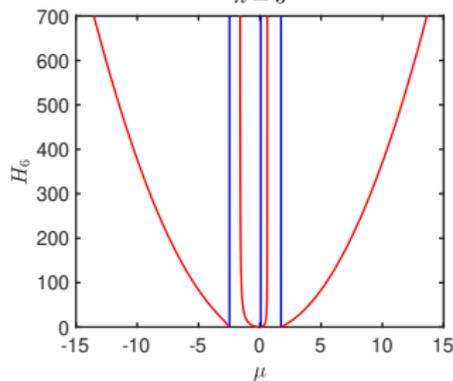
comes from Christoffel–Darboux formula.

Example of the evolution of the eigenvalues with  $H_{2n}$

$$S_3 = -1 \\ n = 2$$



$$(S_3, S_4, S_5) = (-1, 5, -8) \\ n = 3$$



# Hyperbolicity - Eigenvalues

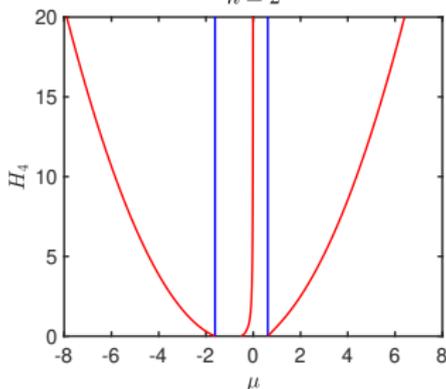
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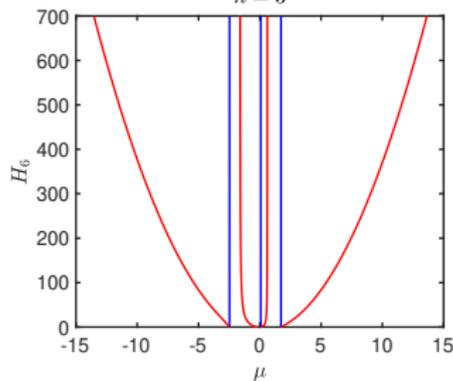
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The moment system with the HyQMOM closure is then hyperbolic, whatever the strictly realizable moment.

# Practical computations

## Closure, directly from the moments $\mathbf{m}_{2n}$

- 1 compute the  $(\bar{a}_k)_{k=0}^{n-1}$  and  $(\bar{b}_k)_{k=0}^n$  from  $\mathbf{m}_{2n}$  with the Chebyshev algorithm
- 2 set the closure  $\bar{a}_n = \frac{1}{n} \sum_{k=0}^{n-1} \bar{a}_k$
- 3 compute  $m_{2n+1}$  using the reverse Chebyshev algorithm





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# The 1D Riemann problem

## problem at the kinetic level

Two homogeneous sprays, with Gaussian distribution and infinite Stokes, crossing.

Problem at the kinetic level

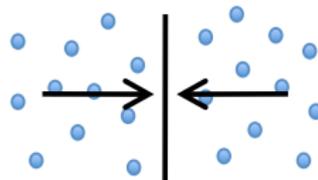
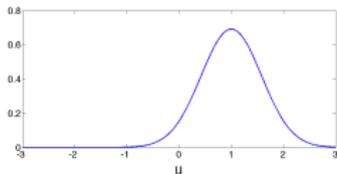
$$\begin{aligned}\partial_t f + \partial_x(vf) &= 0, \\ f(v; 0, x) &= \mathcal{M}_\sigma(v - \bar{u}(x))\end{aligned}$$

with  $\sigma = 1/3$

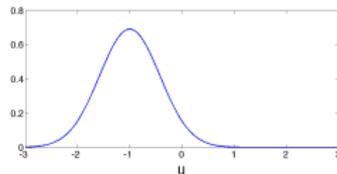
$$\bar{u}(x) = \begin{cases} 1 & \text{if } x < 0, \\ -1 & \text{otherwise.} \end{cases}$$

Analytical solution  $f(t, x, v) = \mathcal{M}_\sigma(v - \bar{u}(x - vt)) = \begin{cases} \mathcal{M}_\sigma(v - 1) & \text{if } v > x/t, \\ \mathcal{M}_\sigma(v + 1) & \text{otherwise.} \end{cases}$

$t = 0$ :



$X=0$



# The 1D Riemann problem

## problem at the kinetic level

Two homogeneous sprays, with Gaussian distribution and infinite Stokes, crossing.

Problem at the kinetic level

$$\begin{aligned}\partial_t f + \partial_x(vf) &= 0, \\ f(v; 0, x) &= \mathcal{M}_\sigma(v - \bar{u}(x))\end{aligned}$$

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## moment problem

$$\partial_t m_k + \partial_x m_{k+1} = 0, \quad k = 0, \dots, 2n$$

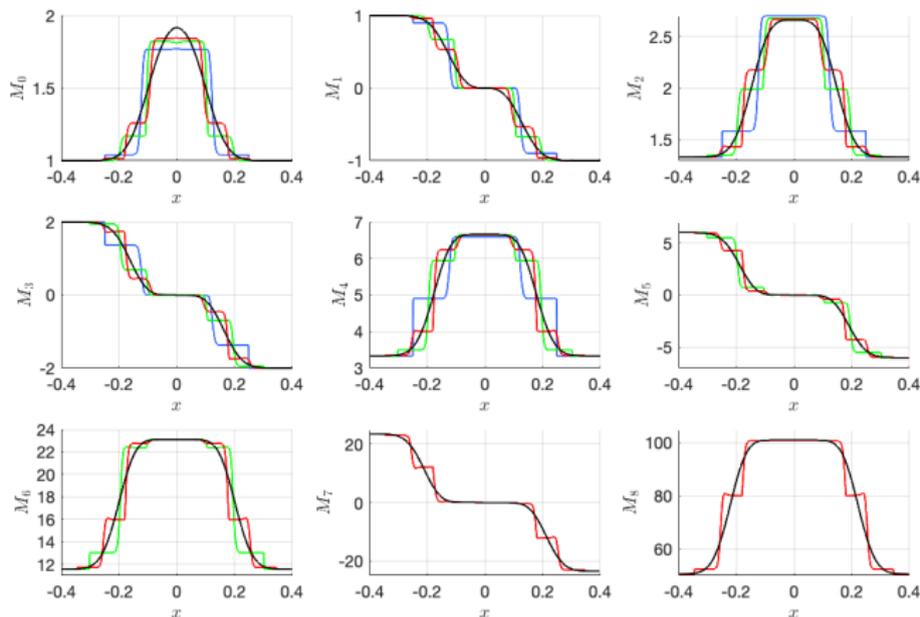
with the initial condition for the standardized moments

$$\rho(0, x) = 1, \quad u(0, x) = \bar{u}(x), \quad C_2(0, x) = \sigma, \quad \begin{cases} S_{2k-1} = 0, \\ S_{2k} = (2k-1)S_{2k-2}, \end{cases} \quad k = 2, \dots, n$$

numerical scheme: HLL [Harten et al., 1983]

## Results

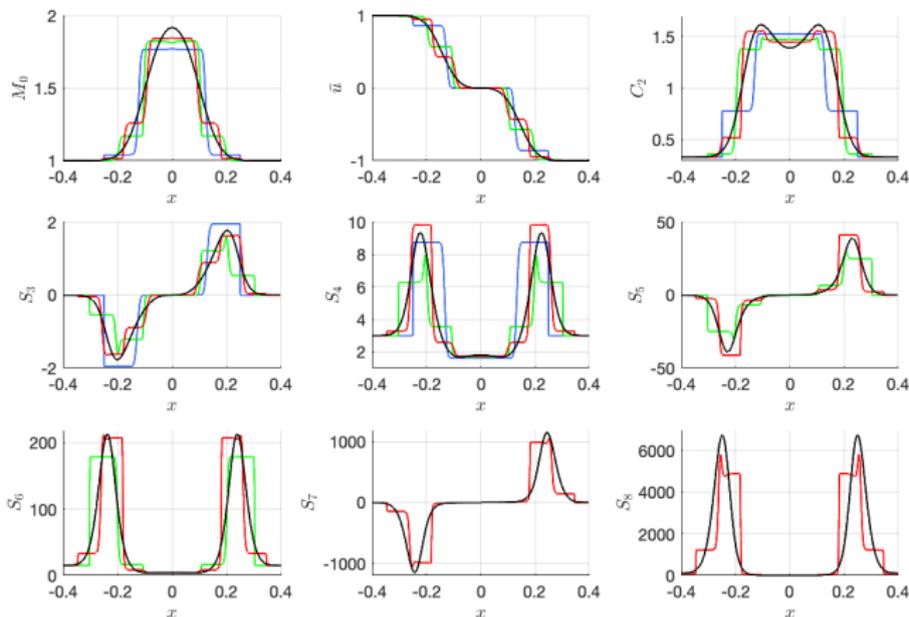
## The 1D Riemann problem - Results

moments - cases  $n=2,3,4$ 

Good behavior on this hard test case.

## Results

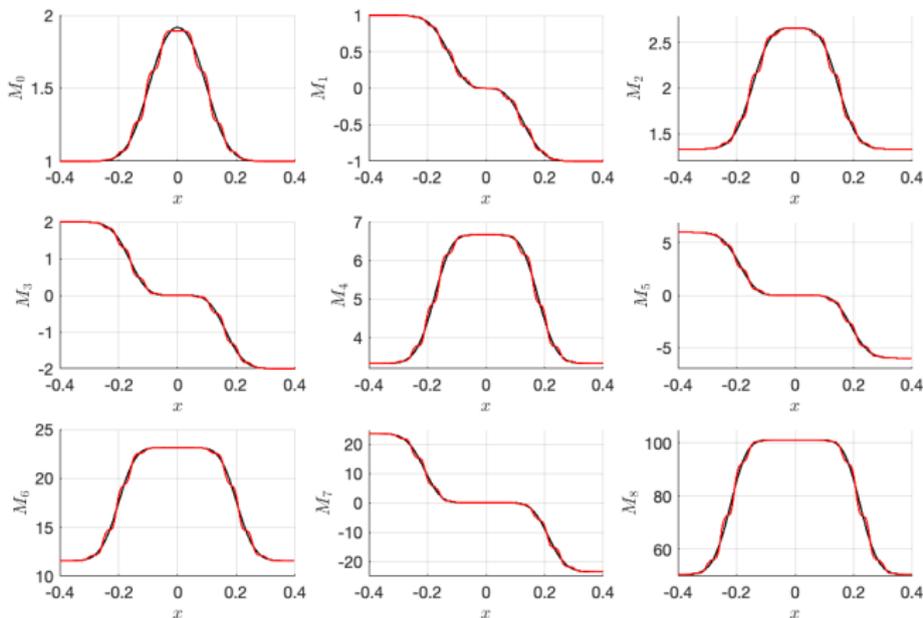
## The 1D Riemann problem - Results

standardized moments - cases  $n=2,3,4$ 

Good behavior on this hard test case.

# The 1D Riemann problem - Results

first moments - case  $n=10$

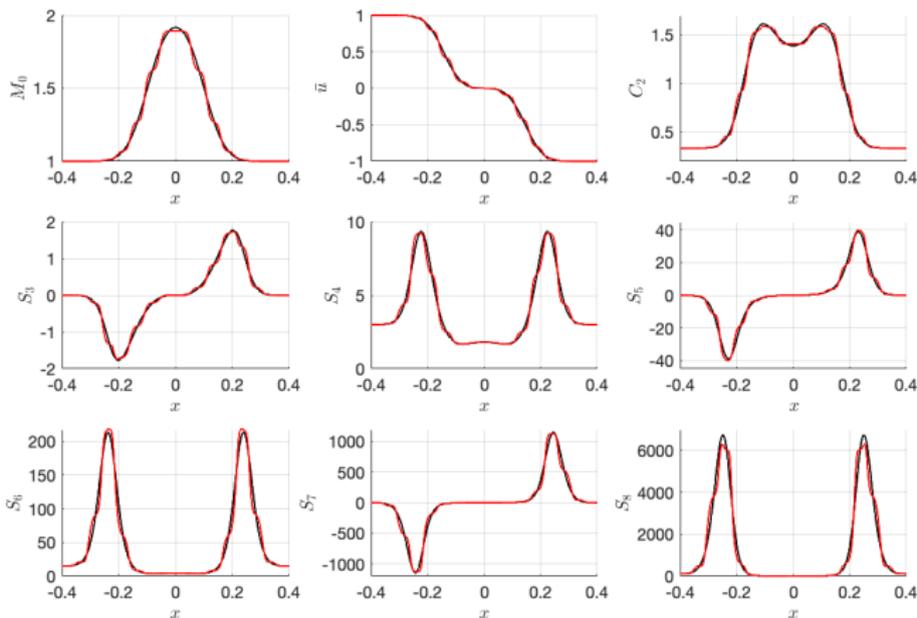


Close to the analytical solution.

## Results

## The 1D Riemann problem - Results

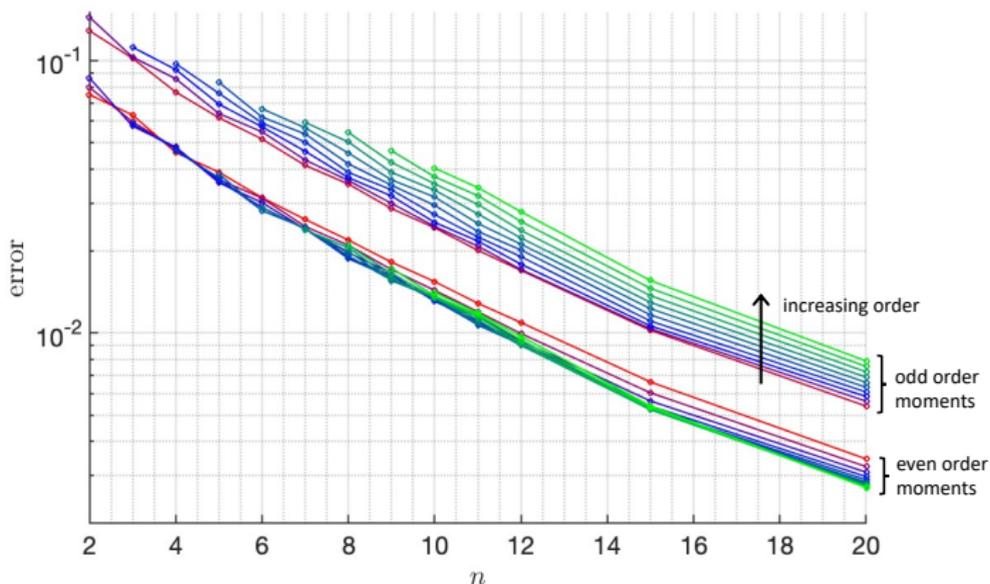
first standardized moments - case n=10



Close to the analytical solution.

# The 1D Riemann problem - Convergence

error on the moments



The moment method seems to converge to the solution of the kinetic equation when the number of moments increases.

# Outline

- 1 Introduction
  - Context
  - Moment method
  - Hyperbolicity
- 2 QMOM
  - Principle of the method
  - Hyperbolicity
- 3 HyQMOM
  - First version of HyQMOM
  - New HyQMOM closure
  - Properties - Practical computations
- 4 Results
  - Configuration
  - Results
- 5 Conclusion, Perspectives

# Conclusion and Perspectives

## Conclusion

- Closure inducing a global hyperbolicity
- Include the Maxwellian distribution
- Good behavior at the boundary of the moment space
- Efficient algorithm to compute the closure and the eigenvalues

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- 2D-3D version of the HyQMOM closure

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## Perspectives

- 2D-3D version of the HyQMOM closure

THANK YOU FOR YOUR ATTENTION

# References I



Chalons, C., Kah, D., and Massot, M. (2012).

Beyond pressureless gas dynamics: quadrature-based velocity moment models.  
*Commun. Math. Sci.*, 10(4):1241–1272.



Chebyshev, P. L. (1859).

Sur l'interpolation par la méthode des moindres carrés.  
*Mém. Acad. Impér. Sci. St. Petersbourg*, 1(15):1–24.  
Also in *œuvres I* pp. 473–498.



Fox, R. O. (2008).

A quadrature-based third-order moment method for dilute gas–particle flow.  
*J. Comput. Phys.*, 227(12):6313–6350.



Fox, R. O. and Laurent, F. (2021).

Hyperbolic quadrature method of moments for the one-dimensional kinetic equation.  
submitted, <https://hal.archives-ouvertes.fr/hal-03171566/>.



Fox, R. O., Laurent, F., and Vié, A. (2018).

Conditional hyperbolic quadrature method of moments for kinetic equations.  
*J. Comput. Phys.*, 365:269–293.



Gautschi, W. (2004).

*Orthogonal Polynomials: Computation and Approximation*.  
Oxford University Press, Oxford, UK.

# References II



Grad, H. (1949).

On the kinetic theory of rarefied gases.  
*Commun. Pure Appl. Math.*, 2(4):331–407.



Harten, A., Lax, P. D., and van Leer, B. (1983).

On upstream differencing and Godunov-type schemes for hyperbolic conservation laws.  
*SIAM Review*, 25(1):35–61.



Huang, Q., Li, S., and Yong, W.-A. (2020).

Stability analysis of quadrature-based moment methods for kinetic equations.  
*SIAM J. Appl. Math.*, 80(1):206–231.



Levermore, C. D. (1996).

Moment closure hierarchies for kinetic theories.  
*J. Stat. Phys.*, 83:1021–1065.



McGraw, R. (1997).

Description of aerosol dynamics by the quadrature method of moments.  
*Aerosol Science and Technology*, 27:255–265.



Müller, I. and Ruggeri, T. (1998).

*Rational Extended Thermodynamics*.  
Springer-Verlag, New York.



Schmüdgen, K. (2017).

*The Moment Problem*, volume 277 of *Graduate Texts in Mathematics*.  
Springer, Cham.

# References III



Shohat, J. A. and Tamarkin, J. D. (1943).

*The Problem of Moments.*

American Mathematical Society, 4th edition.



Wheeler, J. C. (1974).

Modified moments and Gaussian quadratures.

*Rocky Mt. J. Math.*, 4:287–296.

# Chebyshev algorithm

**Three terms recurrence relation** for a sequence  $(Q_k)_{k \geq 0}$  of orthogonal polynomials relative to  $\langle \cdot, \cdot \rangle$ :

$$Q_{k+1}(x) = (x - a_k)Q_k(x) - b_k Q_{k-1}(x).$$

Chebyshev algorithm [Chebyshev, 1859, Wheeler, 1974, Gautschi, 2004]

$$Z_{k,p} = \langle Q_k X^p \rangle$$

$$Z_{-1,p} = 0, \quad Z_{0,p} = m_p$$

$$Z_{k+1,p} = Z_{k,p+1} - a_k Z_{k,p} - b_k Z_{k-1,p}.$$

$$b_0 = m_0, \quad a_0 = \frac{m_1}{m_0}, \quad \forall k > 0 \quad b_k = \frac{Z_{k,k}}{Z_{k-1,k-1}}, \quad a_k = \frac{Z_{k,k+1}}{Z_{k,k}} - \frac{Z_{k-1,k}}{Z_{k-1,k-1}},$$