

SMAI 2021

A posteriori Finite-Volume local subcell correction of
high-order discontinuous Galerkin schemes for the
nonlinear shallow-water equations

Ali HAIDAR

Université de Montpellier 2, institut Montpelliérain Alexander Grothendieck (IMAG)

21 juin 2021



UNIVERSITÉ
DE MONTPELLIER



Introduction

■ The nonlinear Shallow-Water (NSW) equations

$$\partial_t \eta + \partial_x q = 0, \quad (1a)$$

$$\partial_t q + \partial_x (u q + \frac{1}{2} g(\eta^2 - 2\eta b)) = -g\eta \partial_x b. \quad (1b)$$

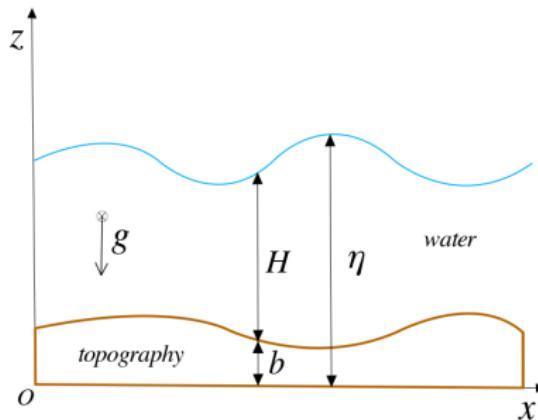


FIGURE – Free surface flow : main notations

Introduction

- Great efforts have been made since the sixties in order to produce accurate approximations of weak solutions of the NSW equations.
- large variety of numerical methods have been developed :
 - **Finite-Volumes (FV) methods**
 - Finite-Elements (FE) methods
 - **Discontinuous Galerkin (DG)**
 - spectral methods
 - residual distribution methods

Introduction : DG methods (advantages)

High-order discontinuous Galerkin (**DG**) methods have become very popular to approximate the solutions of various linear and nonlinear partial differential equations.

- Combines the background of FE and FV methods
- Successfully validated in many domain of applications.
- Reach any Arbitrary high-order of accuracy in space.
- Handle complex geometries.
- Local conservation ; robustness ; strong stability properties.

Introduction : DG methods (disadvantages)

- DG still suffer from the lack of nonlinear stability. In particular, high-order DG methods may produce spurious oscillations in the presence of discontinuities or steeply varying gradients.
- DG may fail near wet/dry zones, negative water height in NSW.

Introduction : FV methods

- **FV advantages** : (1) Low computational cost, (2) their shock-capturing ability, which allows to preserve the discontinuous or steeply varying gradients, (3) and positivity preserving.
- **FV disadvantages** : Low accuracy, low order of convergence, solution dispersion.

Motivation : FV a very good choice for a **local** correction.

Plan du travail

- 1 Discrete formulation**
 - DG formulation
 - DG formulation as a FV scheme on subcells
- 2 *A posteriori* local subcell correction**
 - Admissibility criteria
 - Subcell low-order corrected FV fluxes
 - Well-balancing property
 - Preservation of the water height positivity
- 3 Numerical validation**
 - Time marching algorithm
 - Test cases

Discrete formulation

Let $\Omega \subset \mathbb{R}$ denote the domain (a segment), such that $\Omega = \bigcup_i \omega_i$, where :

- $\omega_i = \left[x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}} \right]$.
- h_{ω_i} is the length of element ω_i .
- $\mathbb{P}^k(\omega_i)$ the space of polynomials in ω_i .

Well-balanced DG scheme

$$\int_{\omega_i} \partial_t v_{h,i} \psi dx = \int_{\omega_i} F(v_{h,i}, b_{h,i}) \partial_x \psi dx - [F\psi]_{i-1/2}^{i+1/2} + \int_{\omega_i} B(v_{h,i}, b_{h,i}) \psi dx, \quad \forall \psi \in \mathbb{P}^k(\omega_i). \quad (2)$$

$$F_{i+1/2} = F(v_{i+1/2}^-, v_{i+1/2}^+, b_{i+1/2}) \quad (3)$$

where $v_{i+1/2}^\pm$ and $b_{i+1/2} = b_{i+1/2}^\pm$ are the interpolated left and right values of ω , interface $x_{i+1/2}$ and F is a consistent numerical flux.

DG formulation as a FV scheme on subcells

$$\partial_t \bar{v}_m^\omega = -\frac{1}{|S_m^\omega|} \left(\hat{F}_{m+\frac{1}{2}}^\omega - \hat{F}_{m-\frac{1}{2}}^\omega \right) + \bar{B}_m^\omega. \quad (4)$$

$$\hat{F}_{m+\frac{1}{2}}^{\omega_i} = F_{\omega_i} \left(\tilde{x}_{m+\frac{1}{2}}^{\omega_i} \right) - C_{m+\frac{1}{2}}^{i-\frac{1}{2}} \left(F_{\omega_i}(x_{i-\frac{1}{2}}) - F_{i-\frac{1}{2}} \right) - C_{m+\frac{1}{2}}^{i+\frac{1}{2}} \left(F_{\omega_i}(x_{i+\frac{1}{2}}) - F_{i+\frac{1}{2}} \right) \quad (5)$$

we refer to [Haidar et al(2021)] and [Vilar(2019)] for more details.

Physical Admissibility Detection (PAD)

We define a sensor function that :

- Check if the sub-mean values \bar{v}_m^{n+1} belongs to Θ .

$$\Theta = \left\{ (H, q) \in \mathbb{R}^2; H \geq 0 \right\}. \quad (6)$$

- Check if there is any *Nan* values.

Subcell Numerical Admissibility Detection (SubNAD)

In order to tackle the issue of spurious oscillations, we enforce a local *Discrete Maximum Principle*, at the subcell level, as follows :

- Check if, for $m = 1, \dots, k+1$, the following inequalities hold :

$$\min(\bar{\eta}_{m-1}^n, \bar{\eta}_m^n, \bar{\eta}_{m+1}^n) \leq \bar{\eta}_m^{n+1} \leq \max(\bar{\eta}_{m-1}^n, \bar{\eta}_m^n, \bar{\eta}_{m+1}^n).$$

The SubNAD criterion relies on a DMP based on subcell mean values, and not the whole polynomial set of values.

Subcell low-order corrected FV fluxes

Replace the subcell mean value \bar{v}_m^{n+1} by a new corrected one $\bar{v}_m^{*,n+1}$, which is computed using a subcell first-order FV scheme of the form :

$$\bar{v}_m^{*,n+1} = \bar{v}_m^n - \frac{\Delta t^n}{|S_m|} \left(F_{m+\frac{1}{2}}^l - F_{m-\frac{1}{2}}^r \right) + \Delta t^n \bar{B}_m, \quad (7)$$

where $F_{m+\frac{1}{2}}^l$, $F_{m-\frac{1}{2}}^r$ are some new subcell *first-order corrected* FV numerical fluxes .

Subcell low-order corrected FV fluxes

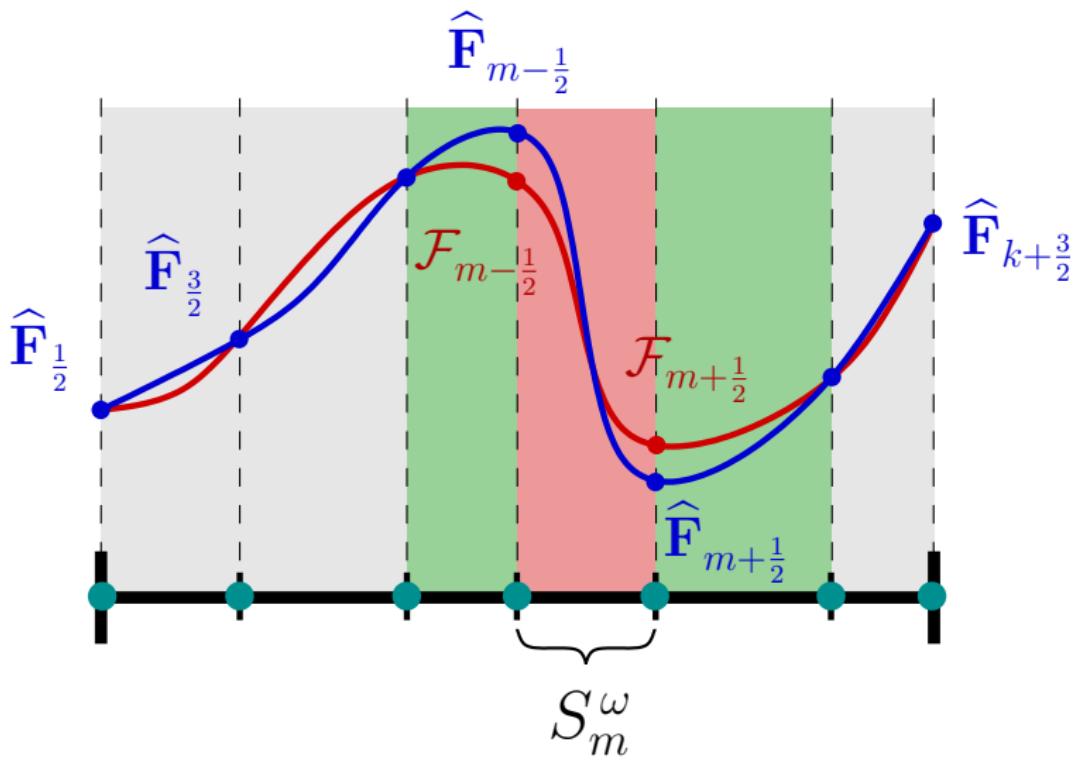
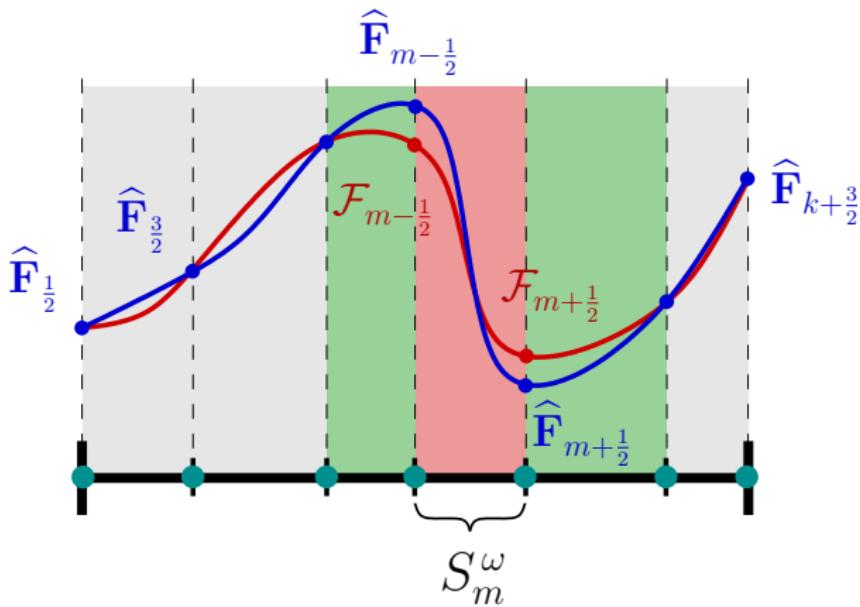


FIGURE – Sketch of the correction of the reconstructed fluxes at subcell boundaries

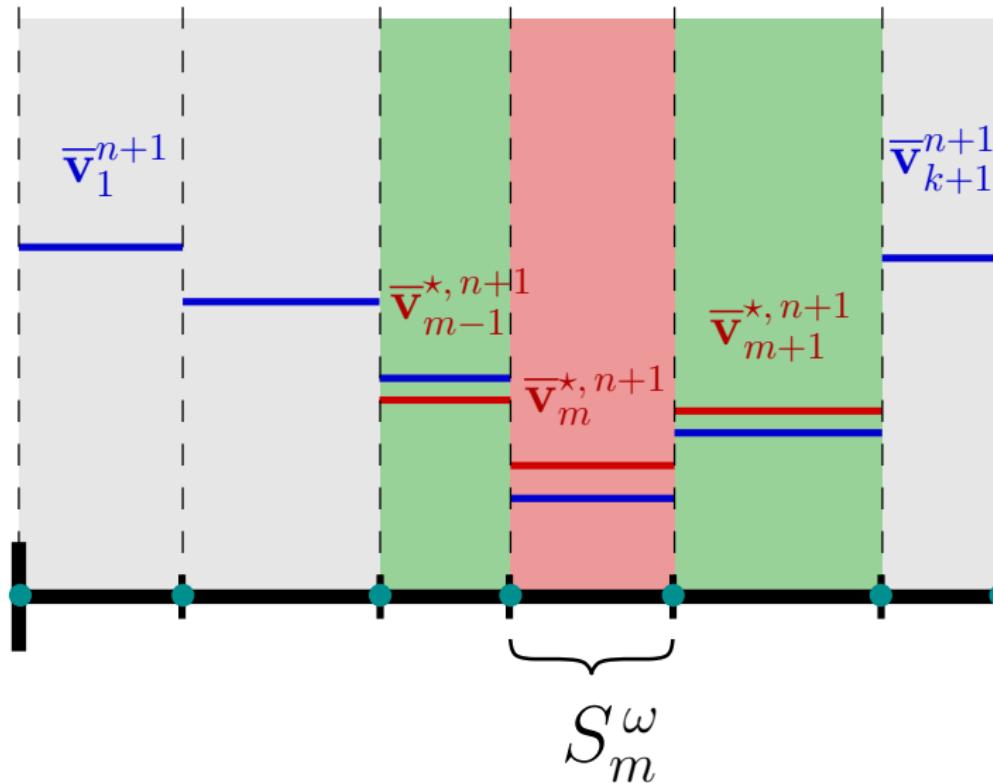
Subcell low-order corrected FV fluxes

$$\bar{v}_{m-1}^{*,n+1} = \bar{v}_{m-1}^n - \frac{\Delta t^n}{|S_{m-1}|} \left(F_{m-\frac{1}{2}}^I - \hat{F}_{m-3/2} \right) + \Delta t^n \bar{B}_{m-1}, \quad (8)$$

$$\bar{v}_{m+1}^{*,n+1} = \bar{v}_{m+1}^n - \frac{\Delta t^n}{|S_{m+1}|} \left(\hat{F}_{m+3/2} - F_{m+\frac{1}{2}}^r \right) + \Delta t^n \bar{B}_{m+1}. \quad (9)$$



Subcell low-order corrected FV fluxes



Well-balancing property (DG numerical flux)

We define the numerical flux function F on interface x as follows :

$$F = F(v_{i+1/2}^-, v_{i+1/2}^+, b_{i+1/2}), \quad (10)$$

where the numerical flux function is the global Lax-Friedrichs flux :

$$F(v^-, v^+, b) = \frac{1}{2} (\mathbf{F}(v^-, b) + \mathbf{F}(v^+, b) - \sigma(v^+ - v^-)), \quad (11)$$

with $\sigma = \max_{\omega \in \Omega} \sigma_\omega$ and

$$\sigma_\omega = \max_m \left(|\bar{u}_m^\omega| + \sqrt{g H_m^\omega} \right).$$

Well-balancing property (FV numerical flux)

We introduce some new Finite-Volume numerical fluxes on subcell's S_m left and right interfaces, denoted by $F_{m-\frac{1}{2}}^r$ and $F_{m+\frac{1}{2}}^l$, as follows :

$$F_{m+\frac{1}{2}}^l = F(\bar{v}_m^+, \bar{v}_{m+1}^-, \bar{b}_m^+) + \begin{pmatrix} 0 \\ g\bar{\eta}_m^+ (\bar{b}_m^+ - b_{\bar{x}_{m+\frac{1}{2}}}^-) \end{pmatrix}, \quad (12)$$

$$F_{m-\frac{1}{2}}^r = F(\bar{v}_{m-1}^+, \bar{v}_m^-, \bar{b}_m^-) + \begin{pmatrix} 0 \\ g\bar{\eta}_m^- (\bar{b}_m^- - b_{\bar{x}_{m-\frac{1}{2}}}^+) \end{pmatrix}. \quad (13)$$

Preservation of the water height positivity

$$\bar{v}_m^{*,n+1} = \bar{v}_m^n - \frac{\Delta t^n}{|S_m|} \left(F_{m+\frac{1}{2}}^l - F_{m-\frac{1}{2}}^r \right) + \Delta t^n \bar{B}_m, \quad (14)$$

This scheme preserve positivity (on a sub-cell scale).

Remark

- Scheme (14) conserve height positivity. Actually we don't need more than that
- We have the well balancing for all schemes, on a sub-cell scale.

Time marching algorithm

We advance in time steps using the explicit SSP-RK schemes. For instance, writing the semi-discrete NSW equation in the operator form

$$\partial_t v_h + A_h(v_h) = 0,$$

we advance from time level n to $(n+1)$ with the third-order scheme as follows:

$$v_h^{n,1} = v_h^n - \Delta t^n A_h(v_h^n),$$

$$v_h^{n,2} = \frac{1}{4}(3v_h^n + v_h^{n,1}) - \frac{1}{4}\Delta t^n A(v_h^{n,1}),$$

$$v_h^{n+1} = \frac{1}{3}(v_h^n + 2v_h^{n,2}) - \frac{2}{3}\Delta t^n A_h(v_h^{n,2}).$$

Test cases : Smooth analytic solution

$$u_0(x) = \begin{cases} 1 & \text{if } x \leq 0, \\ e^{-x^{Ns+1}} & \text{elsewhere.} \end{cases}$$

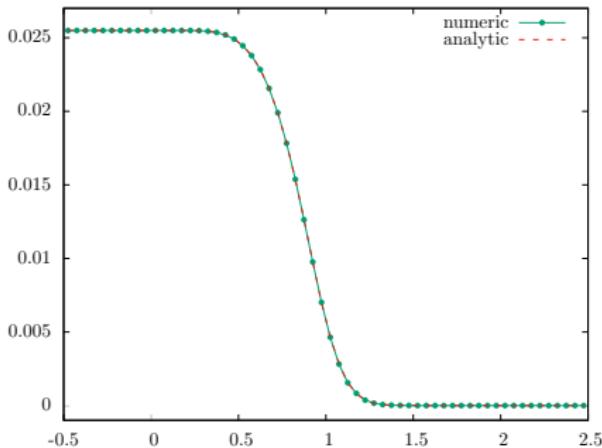


FIGURE – Test 1 - A new analytical solution for the NSW equations - Free surface elevation computed at $t = 0.1$ s with the for $k = 3$ and $n_e = 60$.

Test cases : Smooth analytic solution

k	1		2		3	
h	$E_{L_2}^h$	$q_{L_2}^h$	$E_{L_2}^h$	$q_{L_2}^h$	$E_{L_2}^h$	$q_{L_2}^h$
$\frac{1}{15}$	5.91E-4	1.96	2.13E-5	3.19	3.20E-6	4.05
$\frac{1}{30}$	1.52E-4	2.02	2.33E-6	2.85	1.93E-7	4.18
$\frac{1}{60}$	3.73E-5	2.02	2.99E-7	2.95	1.06E-8	3.95
$\frac{1}{120}$	9.21E-6	-	4.18E-8	-	6.91E-10	-

TABLE – Test 1 - A new analytical solution for the NSW equations : L^2 -errors between numerical and analytical solutions and convergence rates for η at time $t = 0.1\text{s}$

Test cases : discontinuity $t > t_c$

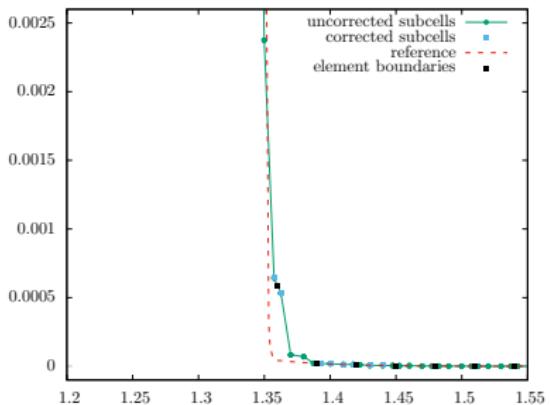
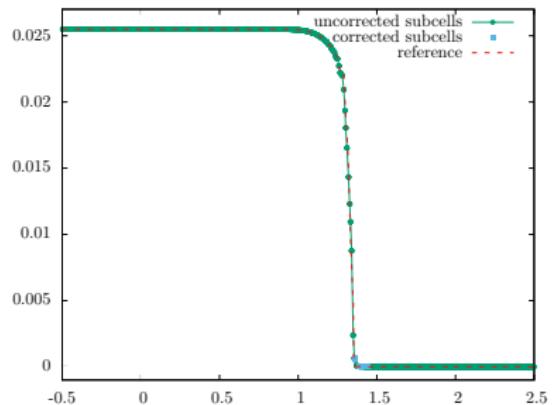
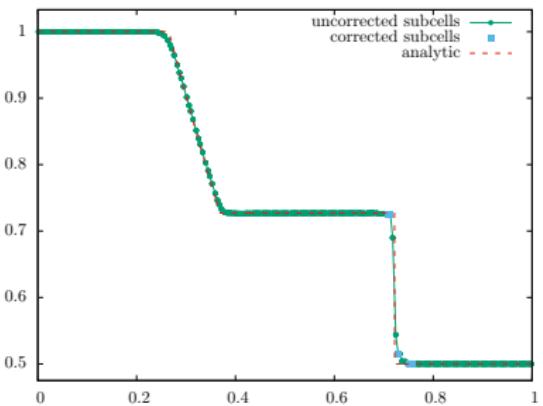
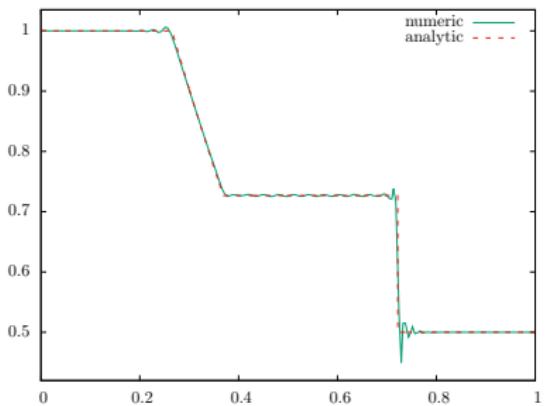
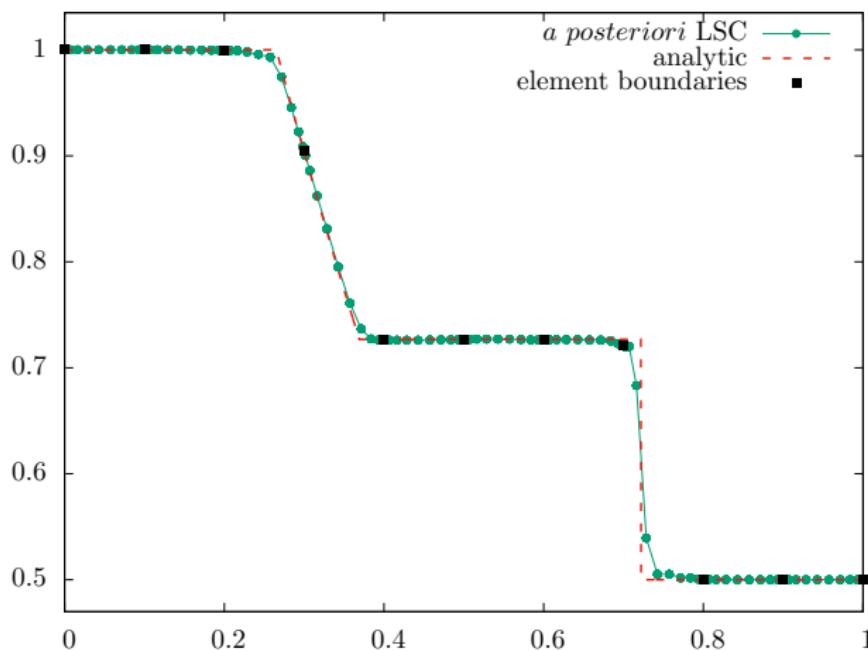


FIGURE – Test 1 - A new analytical solution for the NSW equations, with $k = 3$ and $n_e = 100$, with a zoom on the discontinuity and wet/dry interface.

Test cases : Dam-break (wet bottom)

FIGURE – Test 2 - Dam break on a wet bottom, with $k = 3$ and $n_e = 50$ mesh elements.

Test cases : Dam-break (wet bottom)

FIGURE – Test 2 - Dam break on a wet bottom, with $k = 9$ and $n_e = 10$.

Well-balancing property

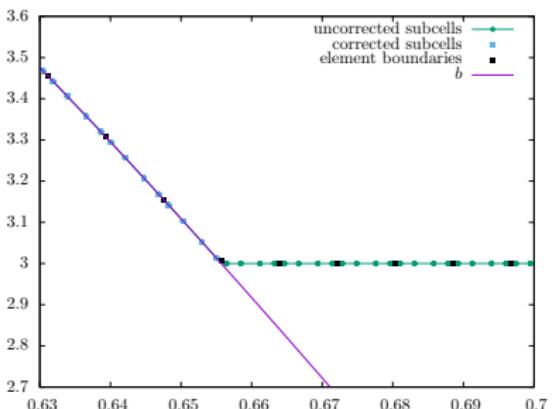
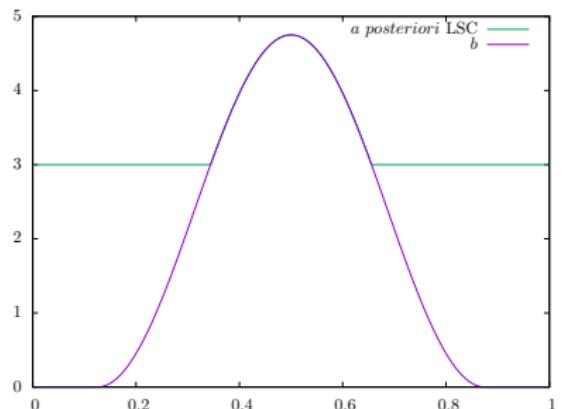


FIGURE – Test 3 - Preservation of a motionless steady state - Free surface elevation at $t = 50s$ (left), with a zoom on the wet/dry interface (right).

Carrier and Greenspan periodic

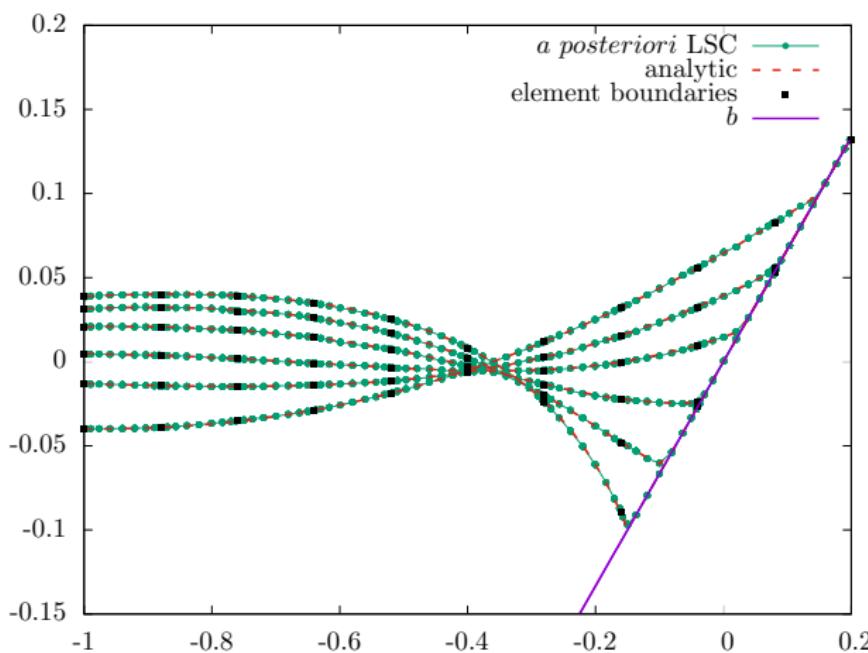


FIGURE – Test 5 - Carrier and Greenspan's periodic solution - Free surface elevation computed for different values of time in the range $[14.5T, 15T]$ for $k = 8$ and $n_e = 10$.

Carrier and Greenspan transient

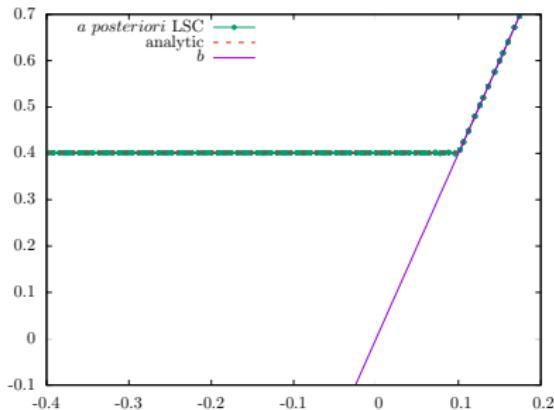
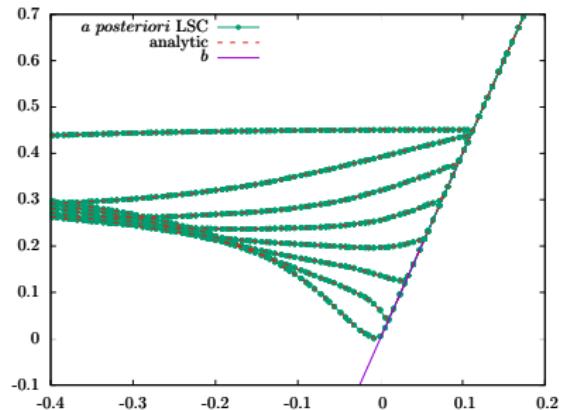


FIGURE – Test 4 - Carrier and Greenspan's transient solution - Free surface elevation for different values of time in the range [0.5 s, 23 s] (left) and at $t = 200$ s (right) for $k = 3$ and $n_e = 50$.

Carrier and Greenspan transient

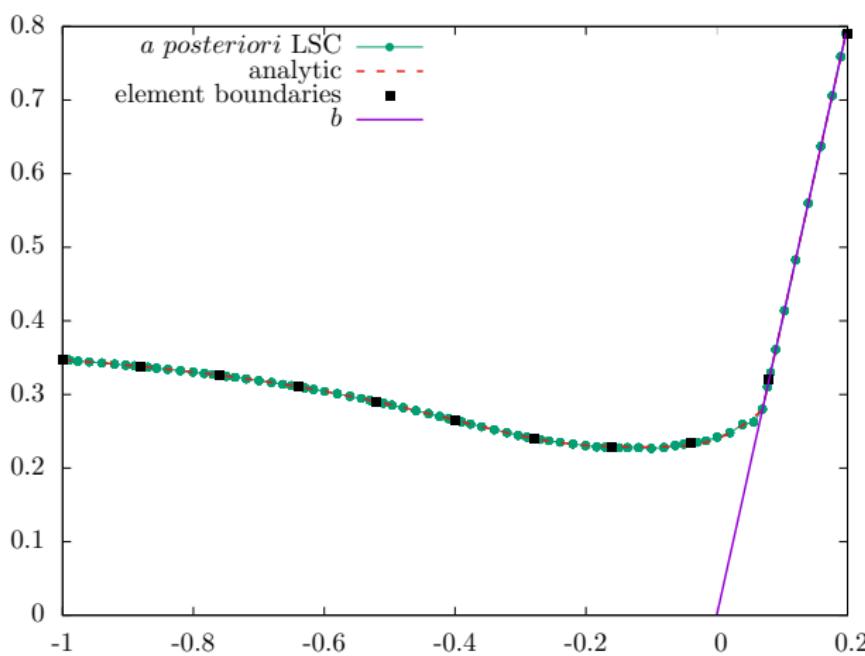


FIGURE – Test 4 - Carrier and Greenspan's transient solution, for $k = 8$ and $n_e = 10$

Conclusion

- New subcell DG-FV formulation for nonlinear Shallow water equation
- The resulting algorithm accurately handles strong shocks with no robustness issues
- It ensures the preservation of the water height positivity at the subcell level
- It preserves the class of motionless steady states (well-balancing)
- It retains the highly accurate subcell resolution of DG schemes
- The procedure is furthermore totally parameter free

Future works

- Investigate moving mesh case, based on an Arbitrary-Lagrangian-Eulerian ALE (in progress)
- Coupling with a floating object.
- Extend to a general 2D case

References I



A. Haidar et al.

A posteriori finite-volume local subcell correction of high-order discontinuous galerkin formulations for the nonlinear shallow-water equations.

J. Comput. Phys., 2021.

Article submitted.



F. Vilar.

A posteriori correction of high-order discontinuous galerkin scheme through subcell finite volume formulation and flux reconstruction.

J. Comput. Phys., 387 :245–279, 2019.