

# Transparent boundary conditions for wave propagation in fractal trees

**Patrick Joly, Maryna Kachanovska**

POEMS (UMR INRIA-CNRS-ENSTA), INRIA, IP Paris

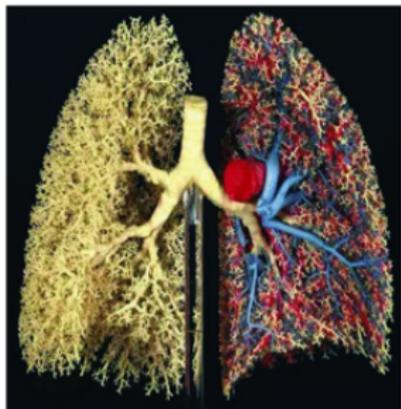
SMAI 2021

# Mathematical Model and Motivation

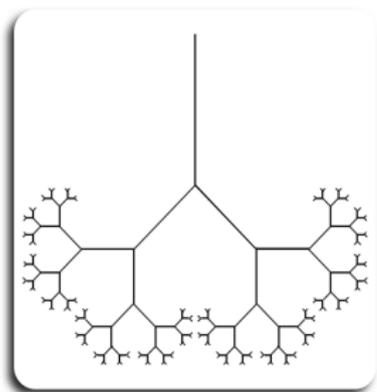
## Wave propagation in lungs

Used in **medical diagnostics**.

Modelled by the **3D wave equation in a thin fractal network** (see works by [B. Maury](#) and co-workers)



**Figure:** Cast of human lungs (photo by Ewald Weibel, University of Bern)



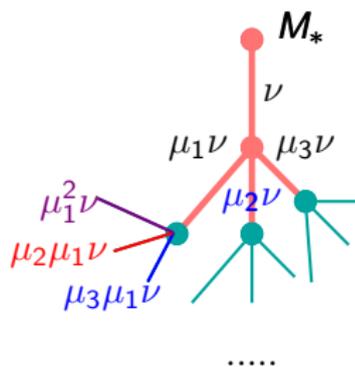
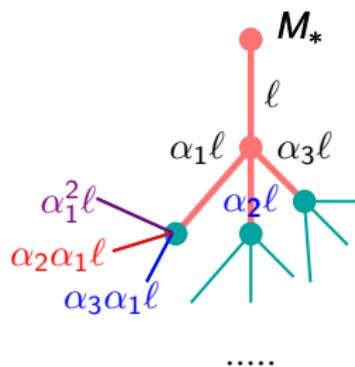
- fractal **infinite 1D** tree  $\mathcal{T}$
- asymptotic analysis of [[Joly, Semin, 2008](#)] when the branch width  $\rightarrow 0$

$$\mu \partial_t^2 u - \partial_s(\mu \partial_s u) = 0$$

# Self-Similarity: Example ( $p$ -adic tree with $p = 3$ )

$\alpha_1, \alpha_2, \alpha_3 < 1$  length ratios

$\mu_1, \mu_2, \mu_3$  weight ratios



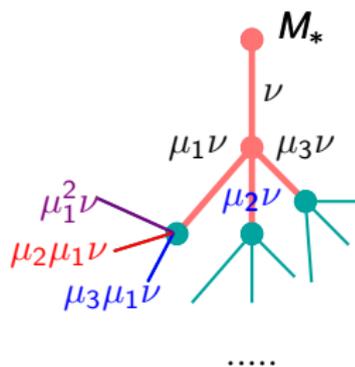
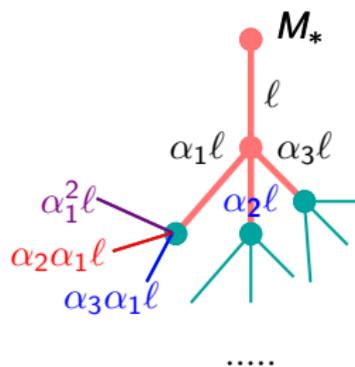
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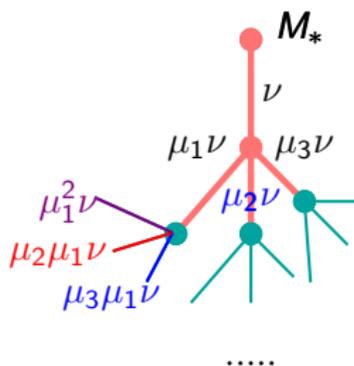
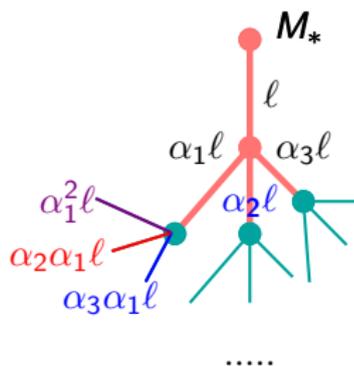
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Lung:  $p = 2$ ,  $\alpha_1 \approx \alpha_2 \approx 0.84$ ,  $\mu_1 \approx \mu_2 \approx 0.75$  (Weibel '63)

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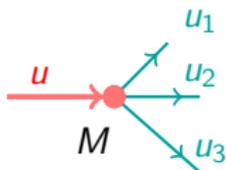
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The model ( $u$  is an acoustic pressure):  $\partial_s(\mu \partial_s u) - \mu \partial_t^2 u = 0$

1D PDE :  $\partial_s^2 u - \partial_t^2 u = 0$  on each branch coupled with

**Continuity:**  $u(M) = u_j(M)$ ,  $j = 1, 2, 3$  in all vertices  $M$

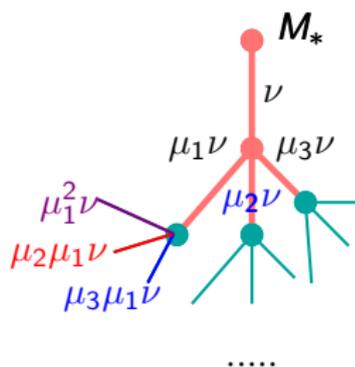
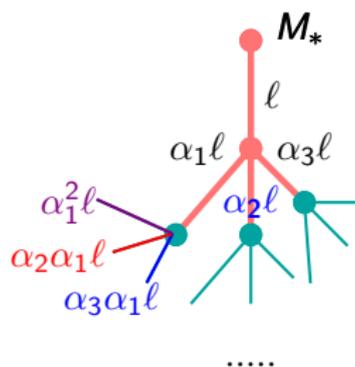
**Kirchoff:**  $\partial_s u(M) = \sum_j \mu_j \partial_s u_j(M)$ , in all vertices  $M$



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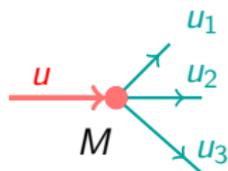
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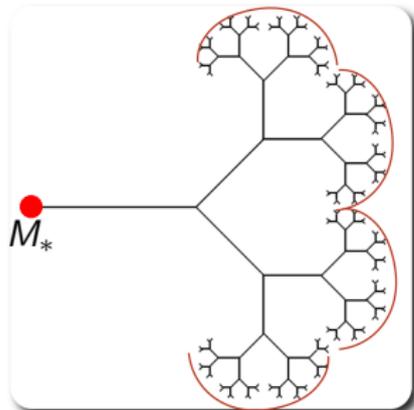
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A trivial example:  $p = 1$ ,  $\mu_1 = 1$ ,  $\alpha_1 < 1$ : a 1D wave equation on an interval

# Boundary Conditions

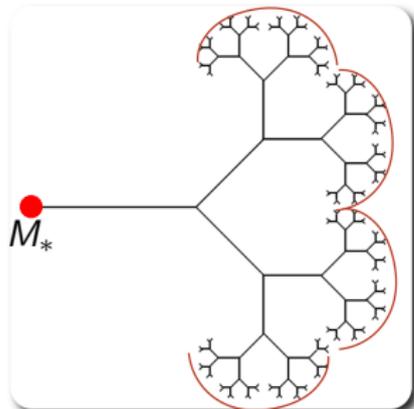


Types of boundary conditions at 'infinity'

We will consider **Neumann BCs**

The BCs will be expressed **variationally**  $\implies$   
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## Variational Framework

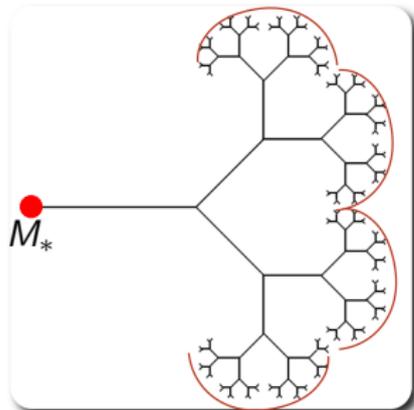
- Square-integrable functions  $L^2_\mu$  on the tree  $\mathcal{T}$ :

$$\|u\|_{L^2_\mu}^2 = \|u\|^2 = \int_{\mathcal{T}} \mu |u|^2 = \sum_{\Sigma \in \mathcal{T}} \int_{\Sigma} \mu_\Sigma |u(s)|^2 ds < \infty.$$

- Sobolev space  $H^1_\mu$ :  $u \in C(\mathcal{T})$ , s.t.

$$\|u\|_{H^1_\mu}^2 = \|u\|_{L^2_\mu}^2 + \int_{\mathcal{T}} \mu |\partial_s u|^2 < \infty.$$

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Neumann (mixed) problem

$$V := \{v \in H_{\mu}^1(\mathcal{T}) : v(M_*) = 0\}$$

$$u \in C^2((0, T]; V) : \int_{\mathcal{T}} \mu \partial_s u \partial_s v + \int_{\mathcal{T}} \mu \partial_t^2 u v = \int_{\mathcal{T}} \mu f v, \quad \text{for all } v \in V$$

+ zero i.c.

## The main objective

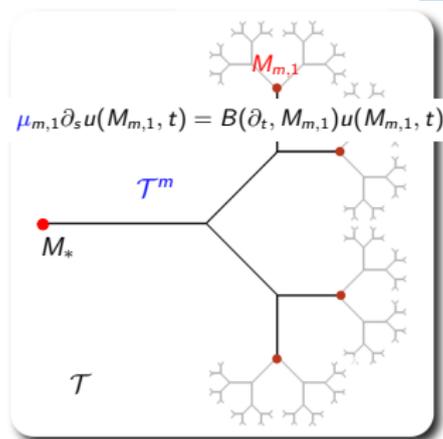
The tree is infinite  $\implies$  restrict computations to a structurally finite domain

Outline:

- 1 transparent boundary conditions for the infinite tree and reference DtN (PhD of A. Semin (2010), P. Joly, MK, A. Semin, *Netw. and Heter. Media* 2019)
- 2 two methods for their approximation
  - convolution quadrature
  - local transparent BCs

# Dirichlet-to-Neumann operators

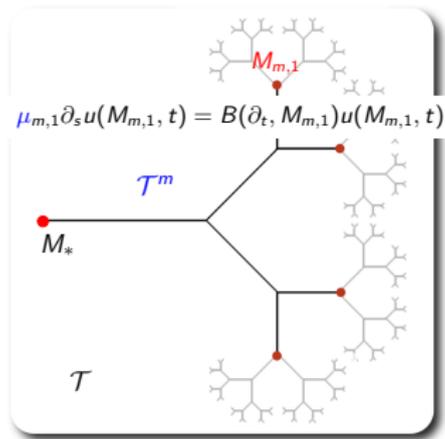
- truncate the tree up to a level  $m \implies$  truncated tree  $\mathcal{T}^m$
- impose a transparent BC  $\mu \partial_s u = B(\partial_t)u$  at the boundary of  $\mathcal{T}^m$  (excl.  $M_*$ )



$B(\partial_t, M_{m,1}) := B_1(\partial_t)$  is a **full-tree** DtN operator

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Coupled formulation:

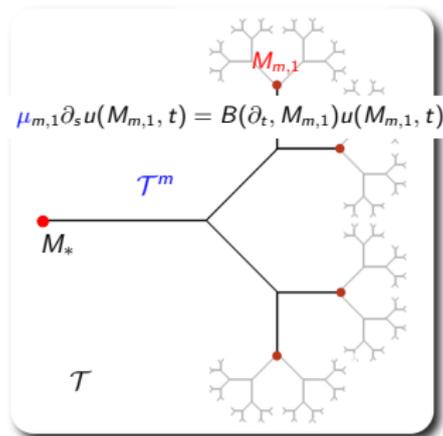
Find  $u \in C^2(0, T; H_\mu^1(\mathcal{T}^m))$ , s.t.  $u(M_*) = 0$  and

$$\int_{\mathcal{T}^m} \mu \partial_t^2 u v + \int_{\mathcal{T}^m} \mu \partial_s u \partial_s v - \sum_{j=1}^{p^m} B(\partial_t, M_{m,j}) u(M_{m,j}, t) v(M_{m,j}) = \int_{\mathcal{T}^m} \mu f v,$$

$\forall v \in H_\mu^1(\mathcal{T}^m)$ , with  $v(M_*) = 0$ .

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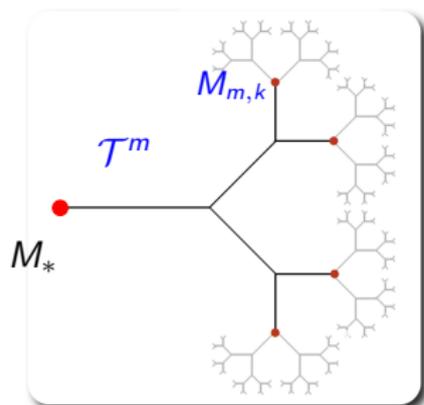
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## Main issues

- the operators  $B(\partial_t)$  are **convolution** operators, **non-local** in time
- no closed form for their kernels or Fourier-Laplace transforms for the kernels

# The reference DtN operator



## Reference DtN:

$\Lambda(\partial_t)g(t): g \rightarrow -\partial_s u(M_*, t)$ , where

$$\mu \partial_t^2 u - \partial_s(\mu \partial_s u) = 0 \text{ on } \mathcal{T},$$

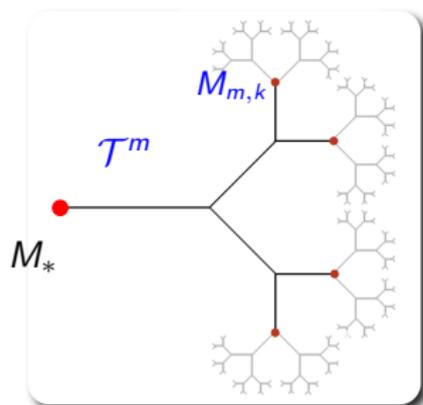
$$u(M_*, t) = g(t),$$

$$u(M_*, 0) = \partial_t u(M_*, 0) = 0, \quad +\text{b.c.}$$

## DtN as a Convolution Operator:

$$\Lambda(\partial_t)g(t) \equiv \int_0^t \lambda(t - \tau)g(\tau)d\tau$$

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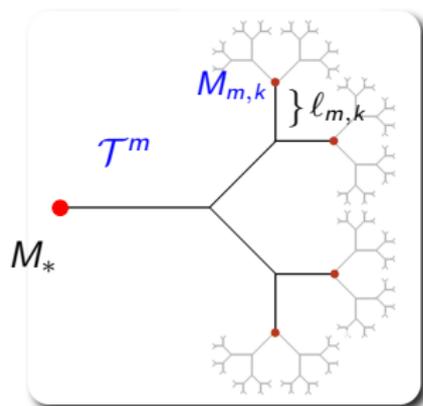
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Fourier-Laplace :

$$(\mathcal{F}g)(\omega) = \int_0^\infty e^{-i\omega t} g(t)dt$$

DtN Symbol :  $\mathbf{\Lambda}(\omega) := \mathcal{F}\lambda(\omega)$

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Using the scaling argument+Kirchoff conditions (see [Joly, MK, Semin '19](#))

$$\mathbf{B}(\omega, M_{m,k}) = - \sum_{k=1}^P \frac{\mu_k \mu_{m,k}}{\alpha_k \ell_{m,k}} \mathbf{\Lambda}(\alpha_k \ell_{m,k} \omega).$$

Conclusion: it suffices to be able to approximate the reference DtN operator.

# The reference DtN operator: Characterization 1

## Compact embedding

The embedding of  $H_{\mu}^1$  into  $L_{\mu}^2$  is **compact**.

Remark: we consider the case  $\alpha_i < 1$  for all  $i$ .

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This means that the spectrum of  $-\mu^{-1}\partial_s(\mu\partial_s)$  on  $V := \{v \in H_\mu^1 : v(0) = 0\}$  is discrete:

Eigenvalues:

Corresponding eigenfunctions:

$$0 < \omega_1^2 \leq \omega_2^2 \leq \dots \rightarrow +\infty,$$
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## Corollary

The reference DtN operator  $\Lambda(\omega)$  is a meromorphic function in  $\mathbb{C}$  analytic in  $\mathcal{B}_\varepsilon(0)$ .  
Moreover,

$$\Lambda(\omega) = \Lambda(0) - \sum_{n=0}^{\infty} \frac{(\partial_s\phi_n(0))^2}{\omega_n^2} \frac{\omega^2}{\omega_n^2 - \omega^2}.$$

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**Problem: computing the eigenvalues and eigenfunctions is expensive!**

## The reference DtN operator: Characterization 2

The non-linear equation defining the DtN

$$\mathbf{\Lambda}(\omega) = \frac{\sum_{k=1}^p \frac{\mu_k}{\alpha_k} \mathbf{\Lambda}(\alpha_k \omega) - \omega \tan \omega}{1 + \omega^{-1} \tan \omega \sum_{k=1}^p \frac{\mu_k}{\alpha_k} \mathbf{\Lambda}(\alpha_k \omega)} \quad (\text{E})$$

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Restriction of the set of the solutions

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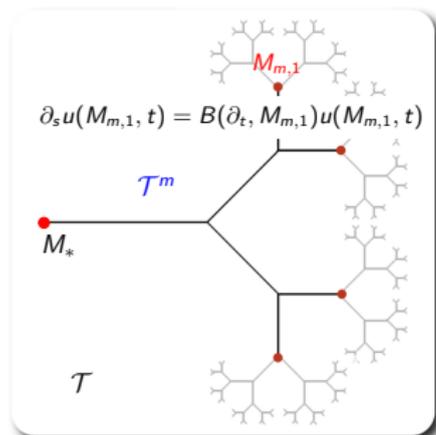
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- knowing  $\Lambda(\alpha_k \omega)$ ,  $\alpha_k < 1 \implies \Lambda(\omega)$
- $\Lambda$  for  $|\omega| < r$  can be found from the truncated **Laurent** series (computable)

# Coupled formulation



$B(\partial_t, M_{m,1}) := B_1(\partial_t)$  is a **full-tree DtN operator**

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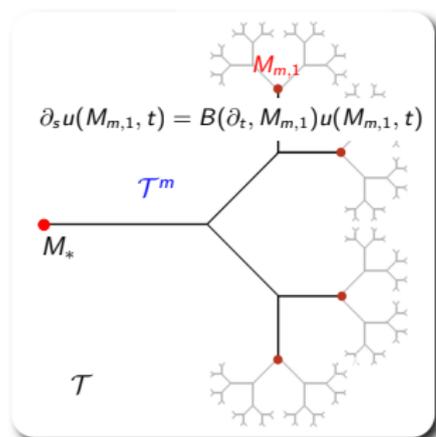
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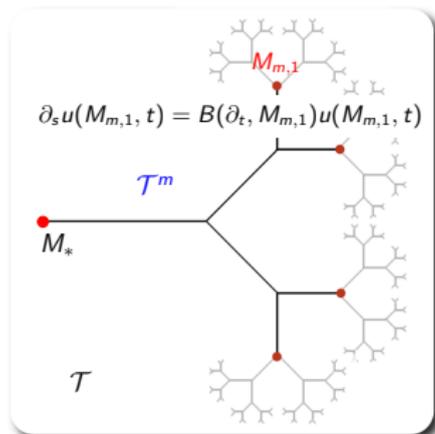
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Computing time-domain convolutions  $B(\partial_t)u$

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- local transparent BCs

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Trapezoid CQ applied to the construction of transparent BCs

- semi-discretize the original problem in time with the  $\theta$ -scheme ( $\theta = \frac{1}{4}$ )

$$\begin{aligned} \mu \frac{u^{n+1} - 2u^n + u^{n-1}}{(\Delta t)^2} - \partial_s(\mu \partial_s \{u^n\}_{1/4}) &= \mu f^n, \\ u^n(M_*) &= 0, \quad u^0 = u^1 = 0, \quad (+\text{b.c.}). \end{aligned} \tag{D}$$

Here  $u^n \sim u(\cdot, n\Delta t)$ ,

$$\{u^n\}_{1/4} = \frac{1}{4} (u^{n+1} + 2u^n + u^{n-1})$$

# Convolution quadrature (CQ)

Ch. Lubich, 1988: discretize  $\int_0^t \mathcal{K}(t-\tau)u(\tau)d\tau$  with multistep (RK) solvers

Trapezoid CQ applied to the construction of transparent BCs

- semi-discretize the original problem in time with the  $\theta$ -scheme ( $\theta = \frac{1}{4}$ )

$$\begin{aligned} \mu \frac{u^{n+1} - 2u^n + u^{n-1}}{(\Delta t)^2} - \partial_s(\mu \partial_s \{u^n\}_{1/4}) &= \mu f^n, \\ u^n(M_*) &= 0, \quad u^0 = u^1 = 0, \quad (+\text{b.c.}). \end{aligned} \tag{D}$$

Here  $u^n \sim u(\cdot, n\Delta t)$ ,

$$\{u^n\}_{1/4} = \frac{1}{4} (u^{n+1} + 2u^n + u^{n-1})$$

- construct exact reference DtN  $\mathbf{\Lambda}_{\Delta t}$ , and the DtNs  $\mathcal{B}_J^{\Delta t}$  for the semi-discrete (D).

# Computing $\Lambda_{\Delta t}$

Continuous problem	Semi-discrete problem	
$-i\omega$ ( $\sim \partial_t$ )	$\underbrace{\frac{2}{\Delta t} \frac{1-z}{1+z}}_{-i\omega + O((\Delta t)^2)}$	$z = e^{i\omega\Delta t}$ Fourier of a shift

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# Coupled formulation

## Implicit form

$$\int_{\mathcal{T}^m} \mu (\partial_t^{\Delta t})^2 \mathbf{u}^n \mathbf{v} + \int_{\mathcal{T}^m} \mu \partial_s \{ \mathbf{u}^n \}_{1/4} \partial_s \mathbf{v} - \sum_j \{ \mathcal{B}_j^{\Delta t} \mathbf{u}^n (M_{m,j}) \}_{1/4} \mathbf{v} (M_{m,j}) = \int_{\mathcal{T}^m} \mu f \mathbf{v}.$$

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## Stability of the coupled formulation

For the continuous formulation the stability is ensured by  $\int_0^t \mathbf{\Lambda}(\partial_t) \mathbf{u} \partial_t \mathbf{u} \geq 0$ .

Discrete formulation:  $\Delta t \sum_{k=0}^n \{ \mathbf{\Lambda}_{\Delta t} \mathbf{u}^k \}_{1/4} \frac{\mathbf{u}^{k+1} - \mathbf{u}^{k-1}}{2\Delta t} \geq 0$

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After the space discretization:

## Explicit form

$$\int_{\mathcal{T}^m} \mu (\partial_t^{\Delta t})^2 \mathbf{u}^n \mathbf{v}_h + \int_{\mathcal{T}^m} \mu \partial_s^h \mathbf{u}_h^n \partial_s^h \mathbf{v}_h - \sum_j \{ \mathcal{B}_j^{\Delta t} \mathbf{u}_h^n (M_{m,j}) \}_{1/4} \mathbf{v}_h (M_{m,j}) = \int_{\mathcal{T}^m} \mu f \mathbf{v}_h.$$

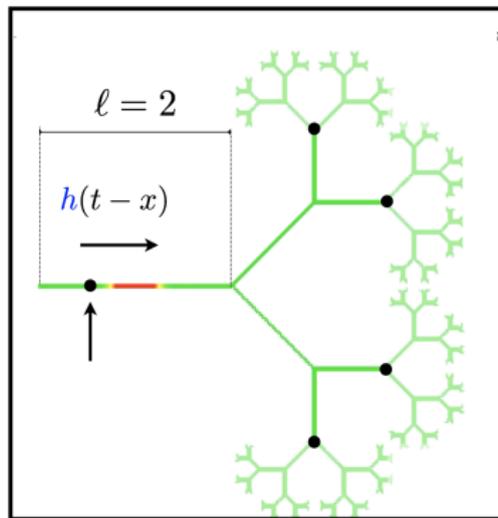
The stability is ensured under the classical CFL condition.

Complexity: evaluating  $\{ \mathcal{B}_m^{\Delta t} \mathbf{u}^n \}_{1/4} \mathbf{v}$  is of  $O(n) \implies O(N^2)$  for  $O(N)$  time steps

# Numerical experiments

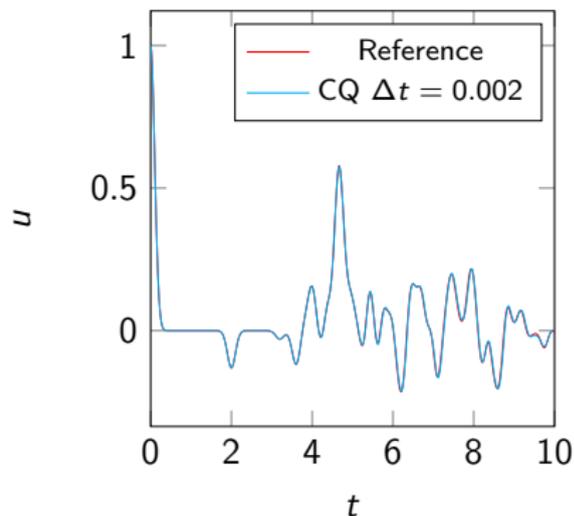
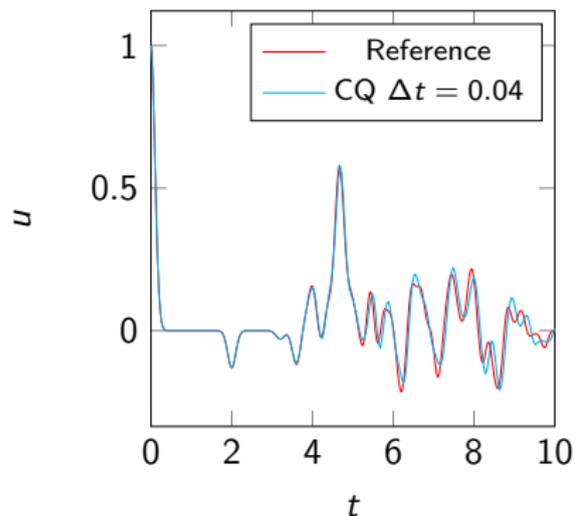
## Tree parameters

$\alpha_1 = 0.4$ ,  $\alpha_2 = 0.3$ ,  $\mu_1 = 1$  and  $\mu_2 = 0.3$ . Neumann case.



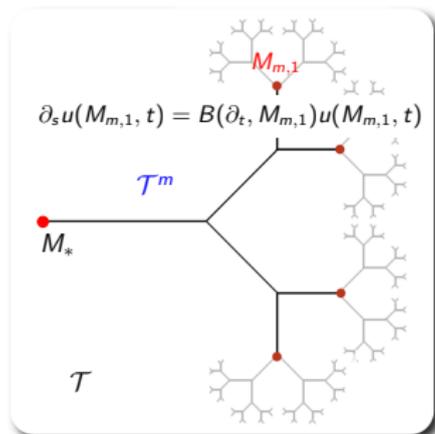
Dirichlet data:  $u(M_*, t) = h(t) = e^{-50(t-1)^2}$ . Cutoff after 3 generations.  
Reference solution: on  $\mathcal{T}^m$  with  $m$  large

# Convolution Quadrature



Remark: Computational times (including pre-computations):  $\sim 1 - 2$  mins for truncated problems;  $\sim 10h$  for reference

# Coupled formulation



$B(\partial_t, M_{m,1}) := B_1(\partial_t)$  is a **full-tree DtN operator**

$$B(\omega, M_{m,k}) = - \sum_{k=1}^p \frac{\mu_k \mu_{m,k}}{\alpha_k \ell_{m,k}} \mathbf{\Lambda}(\alpha_k \ell_{m,k} \omega).$$

**Coupled formulation:**

Find  $u \in C^2(0, T; H_\mu^1(\mathcal{T}^m))$ , s.t.  $u(M_*) = 0$  and

$$\int_{\mathcal{T}^m} \mu \partial_t^2 u v + \int_{\mathcal{T}^m} \mu \partial_s u \partial_s v - \sum_{j=1}^{p^m} B(\partial_t, M_{m,j}) u(M_{m,j}, t) v(M_{m,j}) = \int_{\mathcal{T}^m} \mu f v,$$

$\forall v \in H_\mu^1(\mathcal{T}^m)$ , with  $v(M_*) = 0$ .

Computing time-domain convolutions  $B(\partial_t)u$

- convolution quadrature (Ch. Lubich '88)
- local transparent BCs

# Local transparent BCs

$$\mathbf{\Lambda}(\omega) = \mathbf{\Lambda}(0) - \sum_{n=0}^{\infty} a_n \frac{\omega^2}{\omega_n^2 - \omega^2}, \quad a_n = \frac{(\partial_s \phi_n(0))^2}{\omega_n^2}.$$

Time-domain realization of the reference DtN operator

$$\mathbf{\Lambda}(\partial_t) \mathbf{v} = \mathbf{\Lambda}(0) \mathbf{v} + \sum_{n=0}^{\infty} a_n \frac{d\lambda_n}{dt},$$

$$\frac{d^2}{dt^2} \lambda_n + \omega_n^2 \lambda_n = \frac{d}{dt} \mathbf{v}, \quad \lambda_n(0) = \frac{d}{dt} \lambda_n(0) = 0, \quad 0 \leq n < \infty.$$

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Approximation of the reference DtN operator

$$\mathbf{\Lambda}_N(\partial_t) \mathbf{v} = \mathbf{\Lambda}(0) \mathbf{v} + \sum_{n=0}^{N-1} a_n \frac{d\lambda_n}{dt},$$

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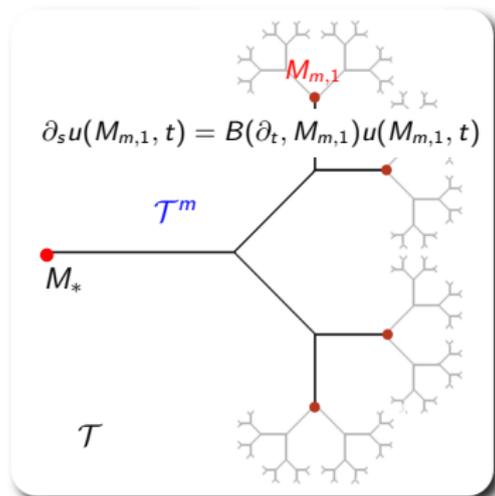
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**Implementation:**  $a_n$  and  $\omega_n^2$  via the residue theorem+contour integration

# Approximating transparent boundary conditions



The 'outer' boundary of  $\mathcal{T}^m$ :  
 $\{M_{m,j}, j = 1, \dots, p^m\}$ .

## Transparent boundary conditions

$$B(\partial_t, M) = - \sum_{k=1}^p \frac{\mu_k \mu_{m,k}}{\alpha_k \ell_{m,k}} \Lambda(\alpha_k \ell_{m,k} \partial_t) \implies$$

$$B_N(\partial_t, M) = - \sum_{k=1}^p \frac{\mu_k \mu_{m,k}}{\alpha_k \ell_{m,k}} \Lambda_N(\alpha_k \ell_{m,k} \partial_t)$$

For brevity:  $B_j(\partial_t) := B(\partial_t, M_{m,j})$ .

## Stability

**Uniform in  $N$  stability:** energy analysis ( $a_n \geq 0$  in  $\Lambda(\omega) = \Lambda(0) - \sum_{n=0}^{\infty} a_n \frac{\omega^2}{\omega_n^2 - \omega^2}$ )

## Error analysis

$u_N$  solution with  $B_N$  on  $\mathcal{T}^m$ ,  $u$  exact. For  $e_N := u_N - u$ :

$$\|\partial_t e_N(\cdot, t)\|_{\mathcal{T}^m} + \|\partial_s e_N(\cdot, t)\|_{\mathcal{T}^m} \leq Ct \|\partial_s u\|_{W^{4,1}(0,t;L^2_\mu)} \sum_{n=N}^{\infty} (\partial_s \phi_n(0))^2 \omega_n^{-4}.$$

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## Goal

Find  $N$  s.t.  $\sum_{n=N}^{\infty} (\partial_s \phi_n(0))^2 \omega_n^{-4} < \varepsilon$

## Reminder

$\omega_n^2$  are eigenvalues of  $-\mu^{-1} \partial_s (\mu \partial_s)$  on  $V := \{v \in H^1_\mu : v(0) = 0\}$ ,  
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## What does not work

Bound  $|\partial_s \phi_n(0)|$  and  $\omega_n$  for every  $n$  (bounding convergent series by a non-convergent one)

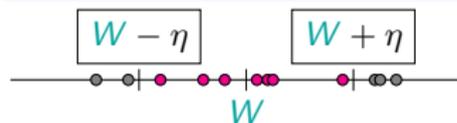
# An alternative idea

## Goal

Find  $N_\varepsilon$  s.t.  $\sum_{n=N_\varepsilon}^{\infty} (\partial_s \phi_n(0))^2 \omega_n^{-4} < \varepsilon$

Adapting [Barnett, Hassell 2011 (Lipschitz domains)]

$$\sum_{\omega_n: |\omega_n - W| \leq \eta} (\partial_s \phi_n(0))^2 < C_\eta W^2.$$



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The estimate does not depend on the number of  $\omega_n$  on the interval  $(W - \eta, W + \eta)$

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Changing the point of view

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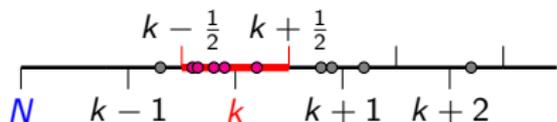
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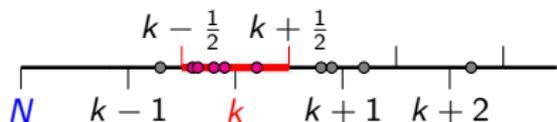
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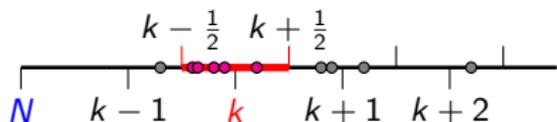
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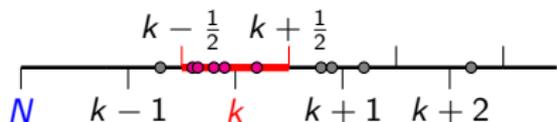
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# An alternative idea

## Goal

Find  $N_\varepsilon$  s.t.  $\sum_{n=N_\varepsilon}^{\infty} (\partial_s \phi_n(0))^2 \omega_n^{-4} < \varepsilon$

Adapting [Barnett, Hassell 2011 (Lipschitz domains)]

$$\sum_{\omega_n: |\omega_n - W| \leq \eta} (\partial_s \phi_n(0))^2 < C_\eta W^2.$$

Changing the point of view

$$\begin{aligned} \sum_{\omega_n \geq N} (\partial_s \phi_n(0))^2 \omega_n^{-4} &\leq \sum_{k=N}^{\infty} \sum_{|\omega_n - k| \leq \frac{1}{2}} (\partial_s \phi_n(0))^2 \omega_n^{-4} \\ &\leq c \sum_{k=N}^{\infty} \frac{1}{k^4} \sum_{|\omega_n - k| \leq \frac{1}{2}} (\partial_s \phi_n(0))^2 \leq c \sum_{k=N}^{\infty} \frac{1}{k^4} k^2 \leq \frac{\tilde{C}}{N}. \end{aligned}$$

Conclusion

$$\sum_{\omega_n \geq \varepsilon^{-1}} (\partial_s \phi_n(0))^2 \omega_n^{-4} < \tilde{C} \varepsilon.$$

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Goal ( $N$  plays a role in the complexity estimate)

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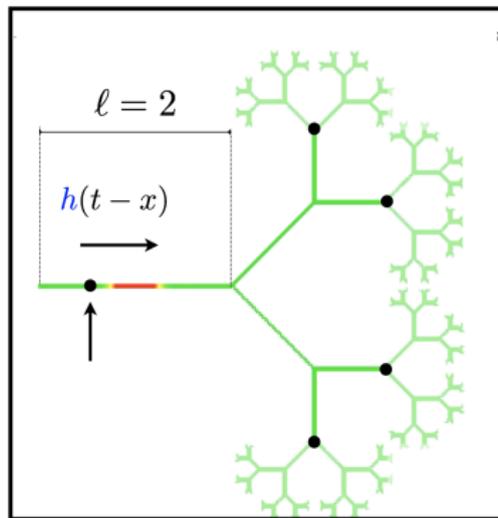
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- $\sum_{k=0}^{p-1} \alpha_k > 1$  ( $d_s > 1$ ):  $N_\varepsilon = O(\varepsilon^{-d_s})$  (Lungs:  $d_s \approx 3.98$ )

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# Numerical experiments

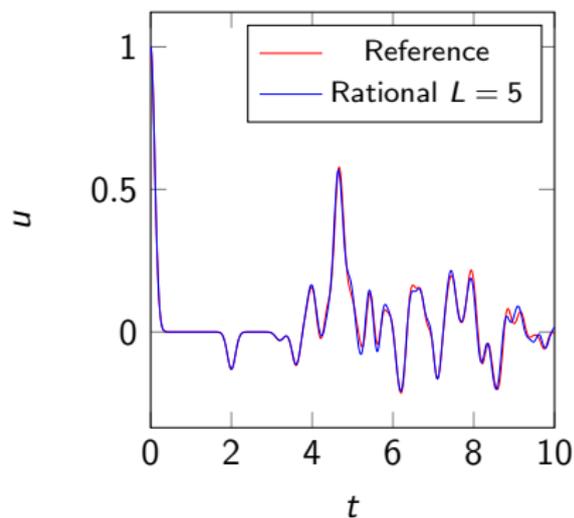
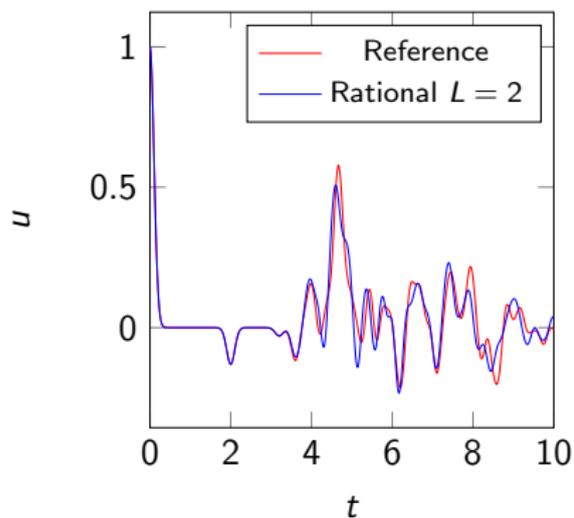
## Tree parameters

$\alpha_1 = 0.4$ ,  $\alpha_2 = 0.3$ ,  $\mu_1 = 1$  and  $\mu_2 = 0.3$ . Neumann case.

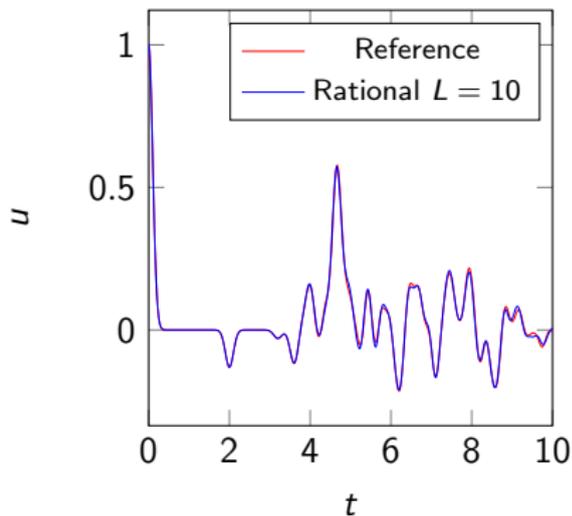
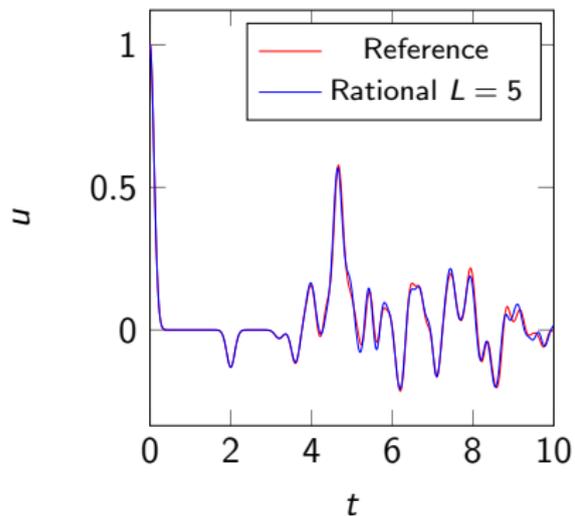


Dirichlet data:  $u(M_*, t) = h(t) = e^{-50(t-1)^2}$ . Cutoff after 3 generations.  
Reference solution: on  $\mathcal{T}^m$  with  $m$  large.

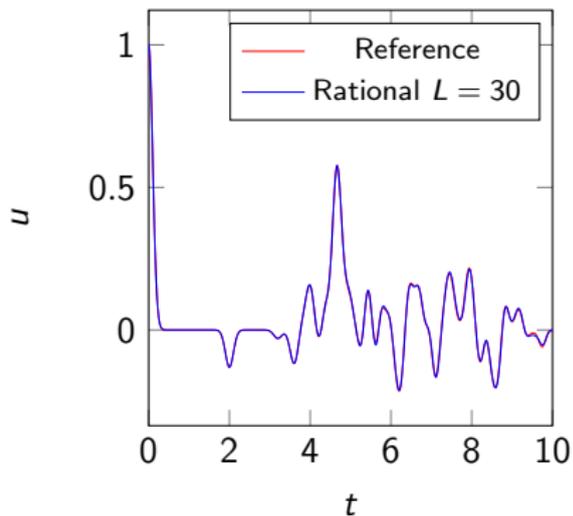
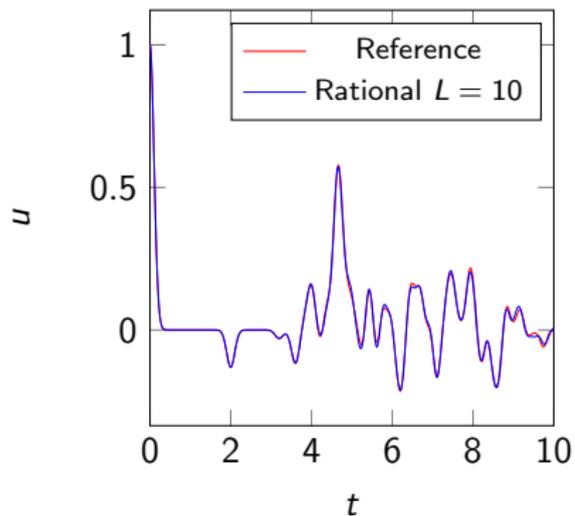
# Rational Function Conditions: Increasing Accuracy



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# Prospectives and future work

- Improved convergence of Local Transparent BCs
- Kirchoff-like conditions accounting for **angles between junctions**
- BCs accounting for the **interaction of the bronchioli with alveoli/lung tissue**

Thank you for your attention!

## References:

- theory and low-order ABC: P. Joly, MK, A. Semin, *Wave Propagation in Fractal Trees. Mathematical and Numerical Issues, Networks and Heterogeneous Media*
- theory and low-order ABC: A. Semin, *Propagation d'ondes dans des jonctions de fentes minces*. Ph.D. thesis (2010)
- Convolution Quadrature: P. Joly, MK, *Transparent Boundary Conditions for Wave Propagation in Fractal Trees: Convolution quadrature*, Num. Math. 2020
- Local Transparent BCs: P. Joly, MK, *Local Transparent Boundary Conditions for Wave Propagation in Fractal Trees, Part (I), (II)*, SISC and SINUM, to appear

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