



On the incompressible limit of a tumor growth model incorporating convective effects

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Introduction

How to model living tissue? A mechanical point of view

- Tissue: **multi-phase fluid**
 - extra cellular matrix
 - proliferating cells
 - dead cells
 - quiescent cells
 - interstitial fluid
 - ...
- Notion of **pressure**:
 - drives the cells **movement**
 - controls the proliferation: **contact inhibition**



Figure 1: Graphical representation of cell division

Macroscopic models

- Systems of PDEs
 - Parabolic - Hyperbolic
(compressible models)
 - Cahn - Hillard models
(4th order)
- Free boundary problems
 - Geometrical or
incompressible models
(Hele Shaw model)
 - Dynamics of $\Omega(t)$

$$\begin{cases} \partial_t n = \nabla \cdot (b(n) \nabla p) \\ p = W'(n) - \delta \Delta n \end{cases}$$

$$\begin{cases} -\Delta p = G(p), \text{ in } \Omega(t) \\ V = -\nabla p \cdot \nu, \text{ on } \partial\Omega(t) \end{cases}$$

Macroscopic models

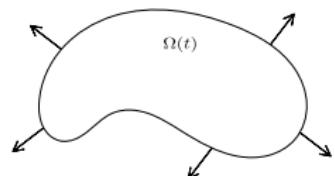
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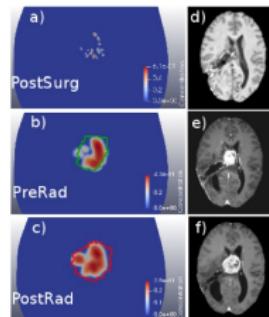
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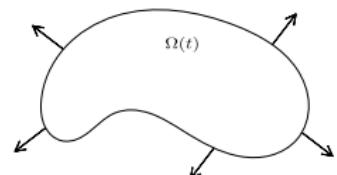
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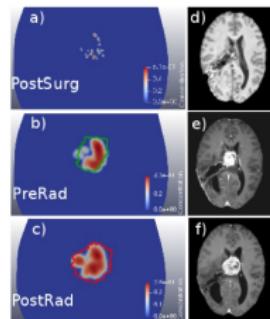
Agosti et al. 2017



Macroscopic models

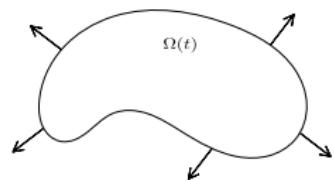
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Agosti et al. 2017

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How can we link *compressible* and *geometrical* models?

Incompressible limit

Mechanical tumor growth model with drift

$$\partial_t n = \underbrace{\nabla \cdot (n \nabla p)}_{\text{pressure effect}} + \underbrace{n G(p)}_{\text{growth term}}$$

- $n(x, t)$ cell population density
- $p(x, t)$ internal pressure
- $\vec{v} = -\nabla p$, Darcy's law

Mechanical tumor growth model with drift

$$\partial_t n = \underbrace{\nabla \cdot (n \nabla p)}_{\text{pressure effect}} - \underbrace{\nabla \cdot (n \nabla \Phi)}_{\text{drift effect}} + \underbrace{n G(p)}_{\text{growth term}}$$

- $n(x, t)$ cell population density
- $p(x, t)$ internal pressure
- $\vec{v} = -\nabla p + \nabla \Phi$,
- $\Phi(x, t)$ concentration of a chemical or a nutrient

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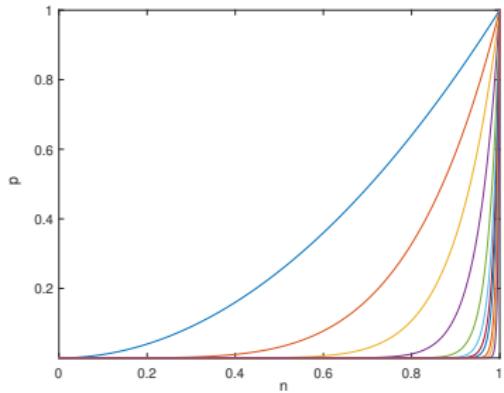
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- $n(x, t)$ cell population density
- $p(x, t)$ internal pressure
- $\vec{v} = -\nabla p + \nabla \Phi$,
- $\Phi(x, t)$ concentration of a chemical or a nutrient
- pressure law of state:

$$p = n^\gamma, \gamma > 1$$

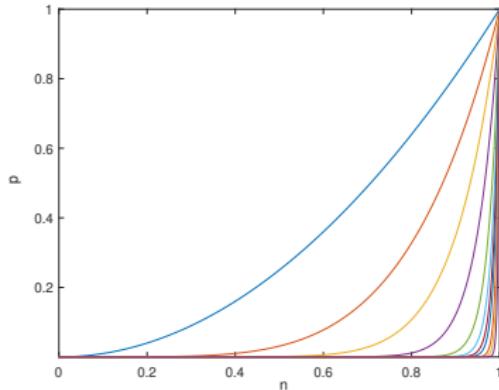
Incompressible limit

$$p = n^\gamma$$



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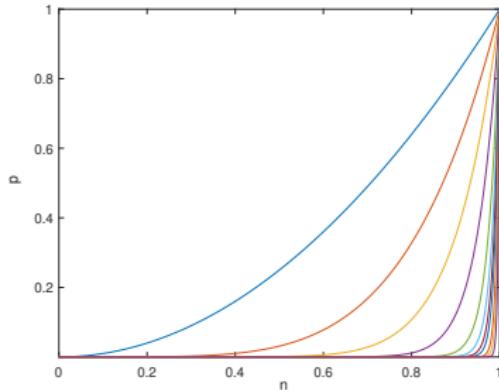


Passing to the limit $\gamma \rightarrow \infty$

$$\begin{cases} p_\infty = 0, & \text{if } n_\infty < 1 \\ p_\infty \in [0, \infty) & \text{if } n_\infty = 1 \end{cases} \Rightarrow p_\infty(1 - n_\infty) = 0$$

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We define $\Omega(t) := \{x; p_\infty(x, t) > 0\} \leftarrow \text{region occupied by the tumor}$

Complementarity relation

Multiplying

$$\partial_t n = \nabla \cdot (n \nabla p) - \nabla \cdot (n \nabla \Phi) + n G(p, c), \quad p = n^\gamma$$

by $\gamma n^{\gamma-1}$

$$\partial_t p = \gamma p (\Delta p - \Delta \Phi + G(p)) + |\nabla p|^2 - \nabla p \cdot \nabla \Phi$$

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Passing, formally, to the limit $\gamma \rightarrow \infty$:

Complementarity relation

$$p_\infty (\Delta p_\infty - \Delta \Phi + G(p_\infty)) = 0$$

Limit model

The **Hele-Shaw problem** reads

$$\begin{cases} -\Delta p_\infty = G(p_\infty) - \Delta \Phi, & \text{in } \Omega(t) = \{x; p_\infty(x, t) > 0\} \\ p_\infty = 0, & \text{on } \partial\Omega(t) \end{cases}$$

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Moreover,

$$\partial_t p_\infty = |\nabla p_\infty|^2 - \nabla p_\infty \cdot \nabla \Phi, \text{ on } \partial\Omega(t)$$

so the normal **boundary velocity** is

$$V = -\partial_\nu p_\infty + \partial_\nu \Phi$$

Free boundary problem

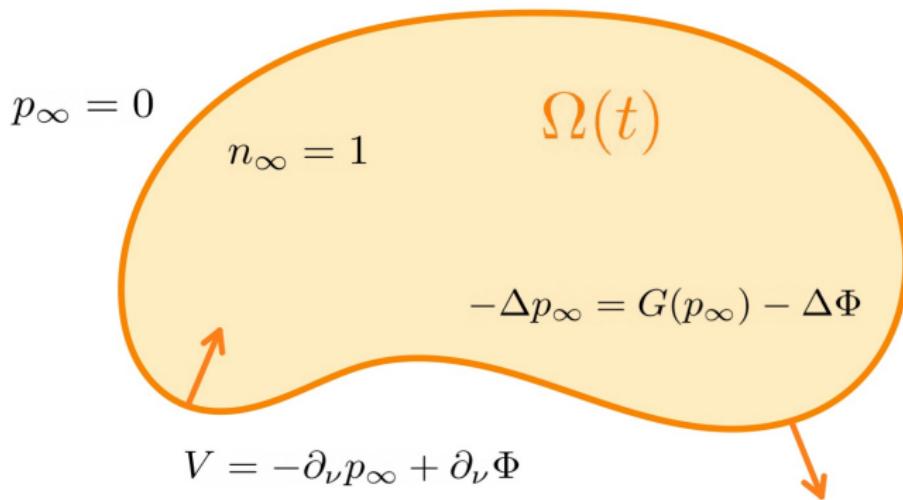


Figure 2: Graphical representation of the limit model when $\gamma \rightarrow \infty$

How to prove it rigorously?

Theorem: limit $\gamma \rightarrow \infty$

$$p_\gamma \rightarrow p_\infty, \quad n_\gamma \rightarrow n_\infty \leq 1 \text{ in } L^1_{x,t}$$
$$\nabla p_\gamma \rightharpoonup \nabla p_\infty \text{ weakly in } L^2_{x,t}$$

$$\partial_t n_\infty = \nabla \cdot (n_\infty \nabla p_\infty) - \nabla \cdot (n_\infty \nabla \Phi) + n_\infty G(p_\infty)$$

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$$p_\infty(\Delta p_\infty - \Delta \Phi + G(p_\infty)) = 0 \quad \text{in} \quad \mathcal{D}'(\mathbb{R}^d \times (0, \infty))$$

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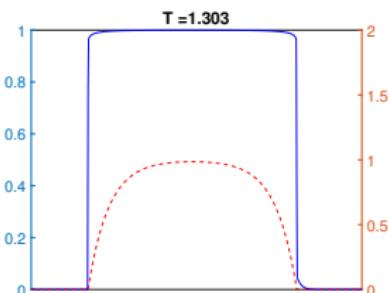
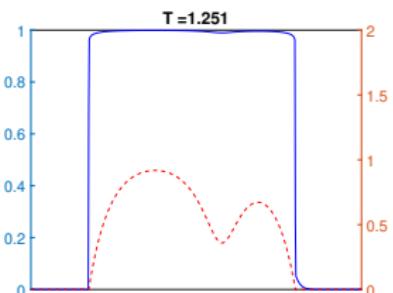
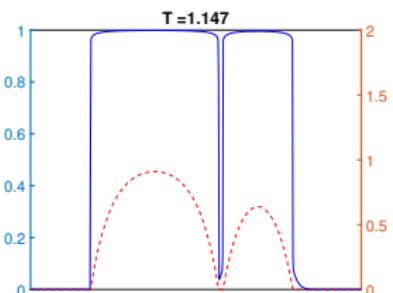
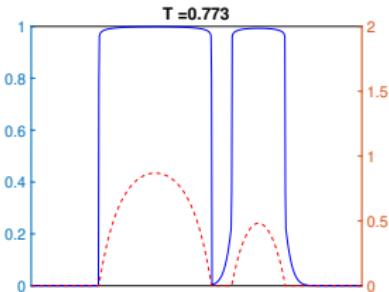
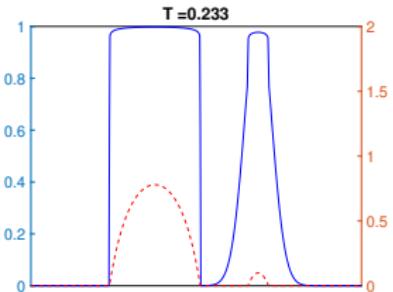
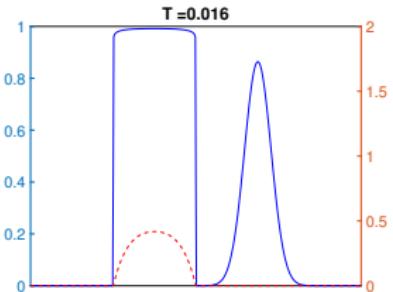
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$$p_\infty(\Delta p_\infty - \Delta \Phi + G(p_\infty)) = 0 \quad \text{in} \quad \mathcal{D}'(\mathbb{R}^d \times (0, \infty))$$

Complementarity relation \iff L^2 -strong compactness of ∇p_γ

Solutions behavior in 1D



Density (blue line), pressure (red dashed line), $\gamma = 90$

Strategy

Uniform a priori estimates:

- $n_\gamma, p_\gamma \in BV(\mathbb{R}^d \times (0, \infty))$
- $\nabla p_\gamma \in L^2(\mathbb{R}^d \times (0, \infty))$

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- an L^4 -optimal bound of ∇p_γ

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- an L^3 -version of the **Aronson-Bénilan estimate**

$$|\Delta p_\gamma + G(p_\gamma)|_- \in L^3(\mathbb{R}^d \times (0, \infty))$$

Important: on the free boundary Δp_γ is a measure: crucial to control Δp_γ itself and not merely $p_\gamma \Delta p_\gamma$

Conclusions and perspectives

Main result

$$p_\infty(\Delta p_\infty - \Delta\Phi + G(p_\infty)) = 0 \quad \text{in} \quad \mathcal{D}'(\mathbb{R}^d \times (0, \infty)).$$

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$$\begin{cases} \partial_t n_1 = \nabla \cdot (n_1 \nabla p) + n_1 F_1(p) + n_2 G_1(p), \\ \partial_t n_2 = \nabla \cdot (n_2 \nabla p) + n_1 F_2(p) + n_2 G_2(p), \\ p = (n_1 + n_2)^\gamma, \quad \gamma > 1 \end{cases}$$

- Incompressible limit? (recent preprint by J.G. Liu and X. Xu)

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- Incompressible limit?
- Existence?

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